

Building Energy Simulation

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Lecture 2: Conduction and Convection

Building Simulation
Introduction
Thermal Analysis
Conduction
Constitutive law
Heat equation
Initial & boundary
Thermal resistance
Thermal circuits
Analysis
1D
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Dynamic
Convection
Constitutive law
Characteristic temp.
Heat rate coefficient
Radiation
Coupled Transfer

Curricula

2 x 4h Lectures

Conduction
Convection
Radiation
Coupled heat transfer

2 x 4h Tutorials and project

Model your own SmartHome
Simulate and discuss

1 x 2h Defend your project

1 x 2h Written exam

Prerequisites

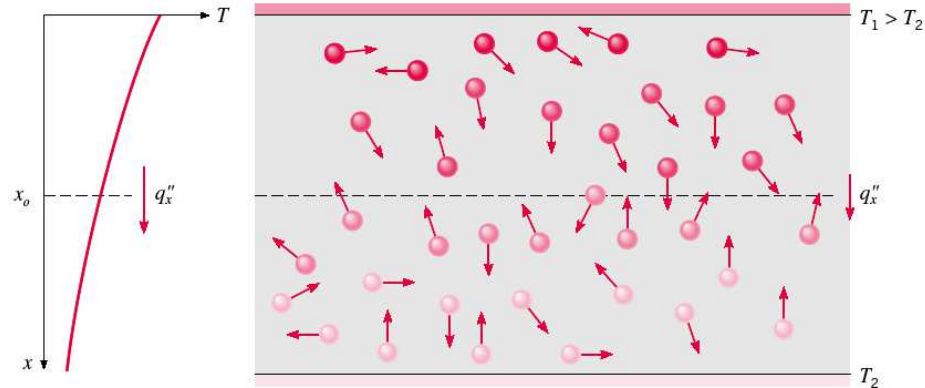
Calculus
Linear algebra
Thermodynamics
Heat and mass trans

Conduction

Constitutive law: Fourier

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• Diffusion



Energy Building Simulation
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Conduction

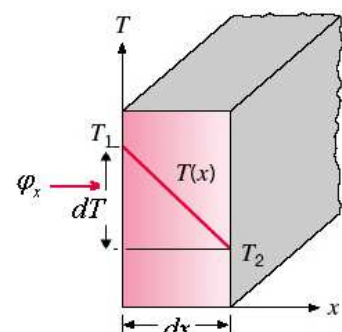
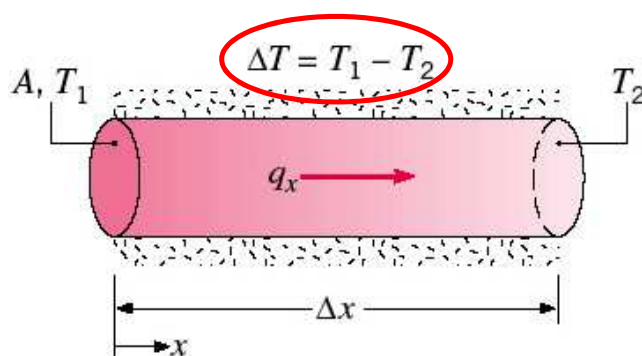
Constitutive law: Fourier

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- Phenomenological law
- Thermal conductivity
 - property of material
 - measured experimentally

$$\dot{Q}_x \equiv q_x = -\lambda A \frac{dT}{dx}$$

$$\varphi = -\lambda \frac{dT}{dx}$$



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Conduction

Constitutive law: Fourier

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- Fourier law:

$$\phi = -\lambda \frac{\partial T}{\partial n} \mathbf{n} \quad \phi = -\lambda \text{grad} T$$

$$\frac{\partial T}{\partial n} \mathbf{n} \equiv \text{grad} T \equiv \nabla T$$

Type de système	Coordonnées	Gradient
Une seule dimension	n	
Trois dimensions		
<i>Cartésiennes</i>	x, y, z	$\text{grad}(T) = \frac{\partial T}{\partial x} \mathbf{x} + \frac{\partial T}{\partial y} \mathbf{y} + \frac{\partial T}{\partial z} \mathbf{z}$
<i>Cylindriques</i>	r, θ, z	$\text{grad}(T) = \frac{\partial T}{\partial r} \mathbf{r} + \frac{1}{r} \frac{\partial T}{\partial \theta} \boldsymbol{\theta} + \frac{\partial T}{\partial z} \mathbf{z}$
<i>Sphériques</i>	r, θ, ω	$\text{grad}(T) = \frac{\partial T}{\partial r} \mathbf{r} + \frac{1}{r} \frac{\partial T}{\partial \theta} \boldsymbol{\theta} + \frac{1}{r \sin \theta} \frac{\partial T}{\partial \omega} \boldsymbol{\omega}$

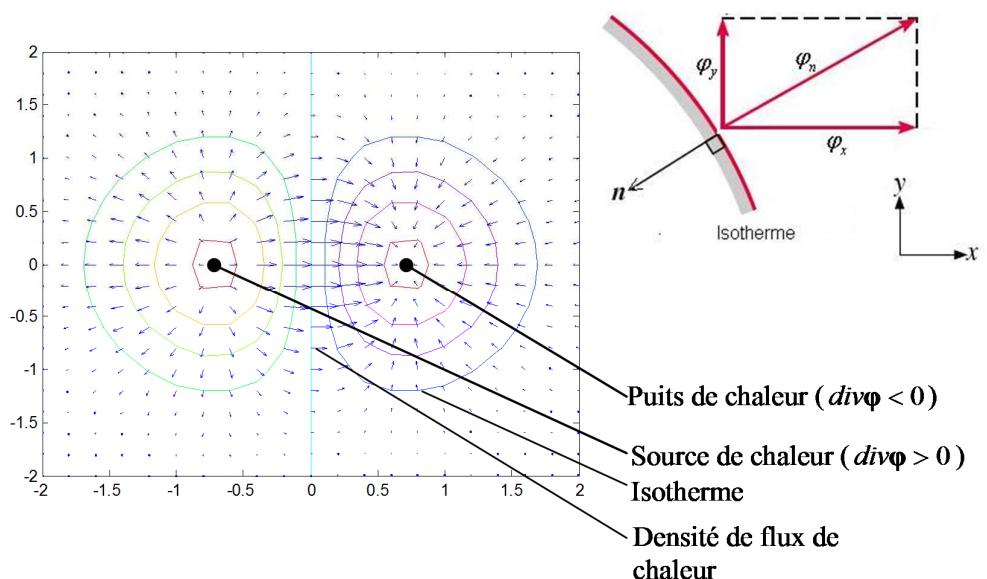
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- temperature distribution: scalar
- heat field: vector



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Heat equation

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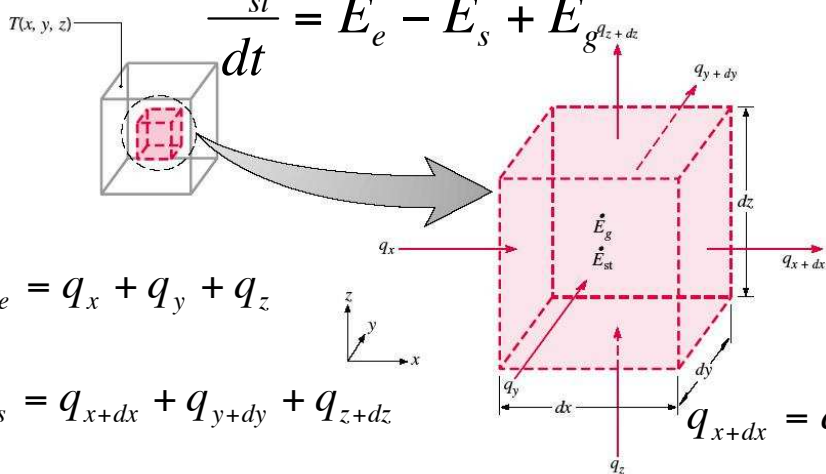
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• Energy balance for a control volume



$$\frac{E_{st}}{dt} = \dot{E}_e - \dot{E}_s + \dot{E}_g$$

$$\dot{E}_e = q_x + q_y + q_z$$

$$\dot{E}_s = q_{x+dx} + q_{y+dy} + q_{z+dz}$$

$$q_{x+dx} = q_x + \frac{\partial q_x}{\partial x} dx$$

$$-\frac{\partial q_x}{\partial x} dx - \frac{\partial q_y}{\partial y} dy - \frac{\partial q_z}{\partial z} dz + p dx dy dz = \rho c dx dy dz \frac{\partial T}{\partial t}$$

$$\underbrace{\dot{E}_e - \dot{E}_s}_{\dot{E}_g} \underbrace{\dot{E}_g} \underbrace{\dot{E}_{st}}$$

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• Energy balance: differential equation

$$\varphi_x \equiv \frac{q_x}{dy dz} \quad \varphi_y \equiv \frac{q_y}{dx dz} \quad \varphi_z \equiv \frac{q_z}{dx dy}$$

$$-\frac{\partial q_x}{\partial x} dx - \frac{\partial q_y}{\partial y} dy - \frac{\partial q_z}{\partial z} dz + p dx dy dz = \rho c dx dy dz \frac{\partial T}{\partial t}$$

$$\underbrace{\dot{E}_e - \dot{E}_s}_{\dot{E}_g} \underbrace{\dot{E}_g} \underbrace{\dot{E}_{st}}$$

$$-div\Phi + p = \rho c \frac{\partial T}{\partial t}$$

$$\frac{\partial T}{\partial t} + div\left(\frac{\Phi}{\rho c}\right) = \frac{p}{\rho c}$$

Conduction

Heat equation

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• Heat equation

$$\frac{\partial T}{\partial t} + \operatorname{div}\left(\frac{\Phi}{\rho c}\right) = \frac{p}{\rho c}$$

continuity equation (fundamental)

$$\Phi = -\lambda \operatorname{grad} T$$

Fourier law (empirical)

$$\rho c \frac{\partial T}{\partial t} = \operatorname{div}(\lambda \operatorname{grad} T) + p$$

heat equation

$$\frac{\partial T}{\partial t} = \operatorname{div}(\alpha \operatorname{grad} T) + \frac{p}{\rho c}$$

Conduction

Heat equation

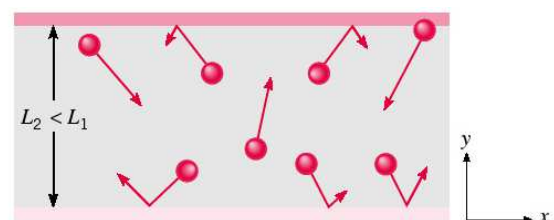
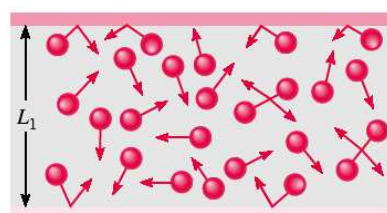
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• Note on empiricism : micro-scale effects

$$\Phi = -\lambda \operatorname{grad} T$$

Fourier law is empirical:

- boundary effects at micro- and nano-scales
- nano-structured materials

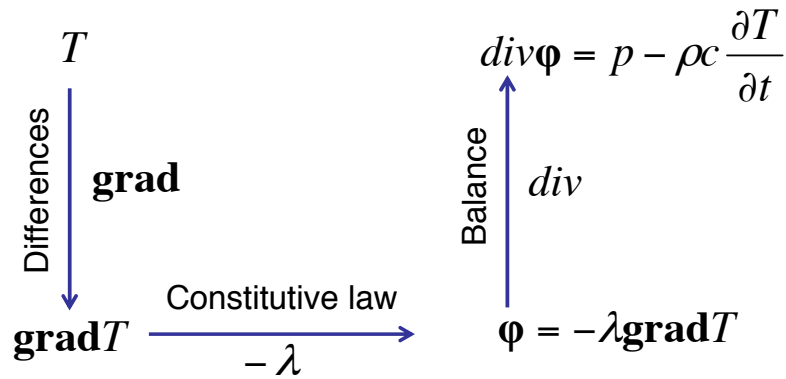


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• Schéma du raisonnement



Symbol	Significance
T	Potential
$\text{grad } T$	Difference of potential
$\phi = -\lambda \text{grad } T$	Heat flow rate due to potential difference
p	External heat flow rates

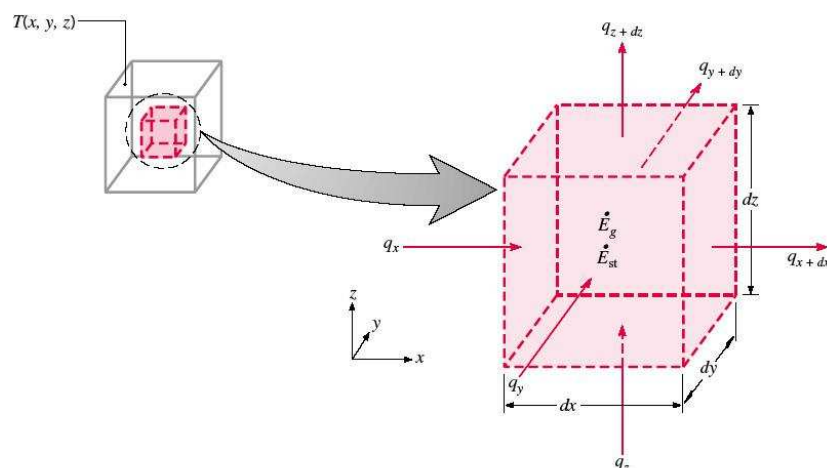
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• Cartesian coordinate system

$$\frac{\partial}{\partial x} \left(\lambda \frac{\partial T}{\partial x} \right) + \frac{\partial}{\partial y} \left(\lambda \frac{\partial T}{\partial y} \right) + \frac{\partial}{\partial z} \left(\lambda \frac{\partial T}{\partial z} \right) + p = \rho c \frac{\partial T}{\partial t}$$



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- Homogeneous and isotropic material

$$\lambda = \text{const.}$$

$$\rho c \frac{\partial T}{\partial t} = \lambda \operatorname{div}(\mathbf{grad} T) + p$$

$$\operatorname{div}(\mathbf{grad} T) = \Delta T \equiv \frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} + \frac{\partial^2 T}{\partial z^2} \quad \text{Laplacien}$$

$$\frac{1}{\alpha} \frac{\partial T}{\partial t} = \operatorname{div}(\mathbf{grad} T) + \frac{p}{\lambda}$$

$$\frac{1}{\alpha} \frac{\partial T}{\partial t} = \Delta T + \frac{p}{\lambda} \quad \alpha = \frac{\lambda}{\rho c} \quad \text{Diffusivity}$$

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- Homogeneous and isotropic material

$$\lambda = \text{const.}$$

in steady-state

$$\frac{\partial T}{\partial t} = 0$$

Poisson equation:

$$\lambda \operatorname{div}(\mathbf{grad} T) + p = 0$$

$$\lambda \Delta T + p = 0$$

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- Homogeneous and isotropic material

$$\lambda = \text{const.}$$

in steady-state

$$\frac{\partial T}{\partial t} = 0$$

without internal sources

$$p = 0$$

Laplace equation:

$$\text{div}(\text{grad}T) = 0$$

$$\Delta T = 0$$

Conduction

Initial and boundary conditions

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$$\rho c \frac{\partial T}{\partial t} = \text{div}(\lambda \text{grad}T) + p$$

For each dimension:

2 conditions at the boundaries

Space distribution of temperature
at initial time

Conduction

Initial and boundary conditions

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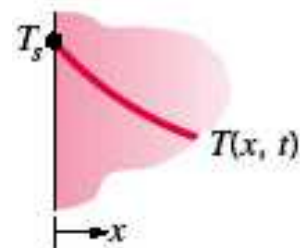
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Initial conditions

$$T|_{t=0} \equiv T_0 = f(x, y, z, 0)$$

Dirichelt boundary conditions

$$T(0, t) = T_s$$



Conduction

Initial and boundary conditions

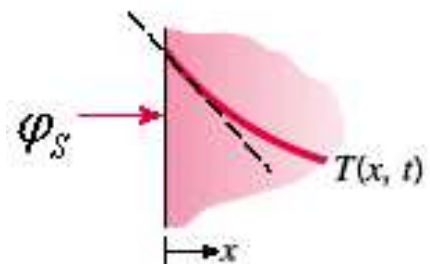
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Neumann boundary condition

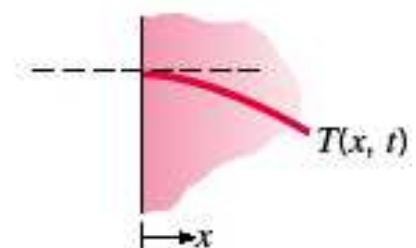
– imposed heat flux

$$-\lambda \frac{\partial T}{\partial x} \Big|_{x=0} = \varphi_s$$



– adiabatic surface or symmetry

$$\frac{\partial T}{\partial x} \Big|_{x=0} = 0$$



Conduction

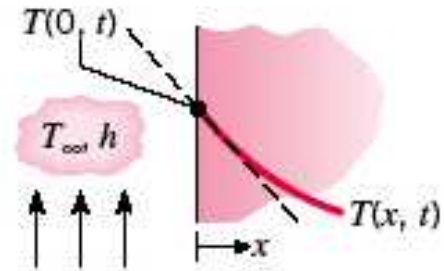
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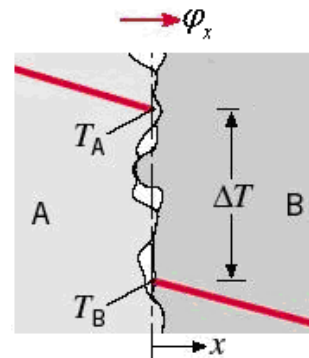
Fourier boundary condition

$$-\lambda \frac{\partial T}{\partial x} \Big|_{x=0} = h_c [T_\infty - T(0, t)]$$



Surfaces in contact

$$\lambda_A \frac{\partial T}{\partial x} \Big|_{x=0^-} = \lambda_B \frac{\partial T}{\partial x} \Big|_{x=0^+}$$



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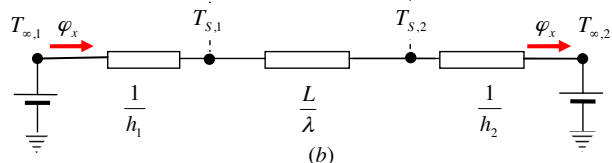
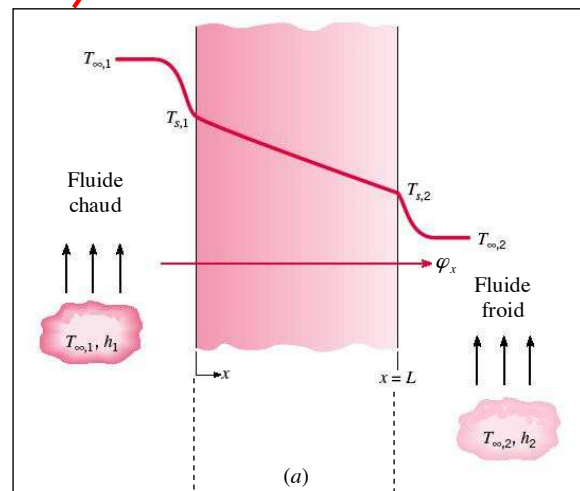
~~$$\rho c \frac{\partial T}{\partial t} = \text{div}(\lambda \text{grad} T) + p$$~~

$$\text{div}(\lambda \text{grad} T) = 0$$

Plane wall

$$\left\{ \begin{array}{l} \frac{d}{dx} \left(\lambda \frac{dT}{dx} \right) = 0 \\ T(0) = T_{s1} \\ T(L) = T_{s2} \end{array} \right.$$

$$T(x) = C_1 x + C_2$$



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• Boundary conditions

$$T(0) = T_{S1} \quad T(L) = T_{S2}$$

→ constants of integration

$$C_2 = T_{S1} \quad C_1 = \frac{T_{S2} - T_{S1}}{L}$$

→ particular solution

$$T(x) = \frac{T_{S2} - T_{S1}}{L} x + T_{S1}$$

• Fourier law →

$$\varphi_x = -\frac{\lambda}{L}(T_{S2} - T_{S1}) = \frac{\lambda}{L}(T_{S1} - T_{S2})$$

• Heat flux

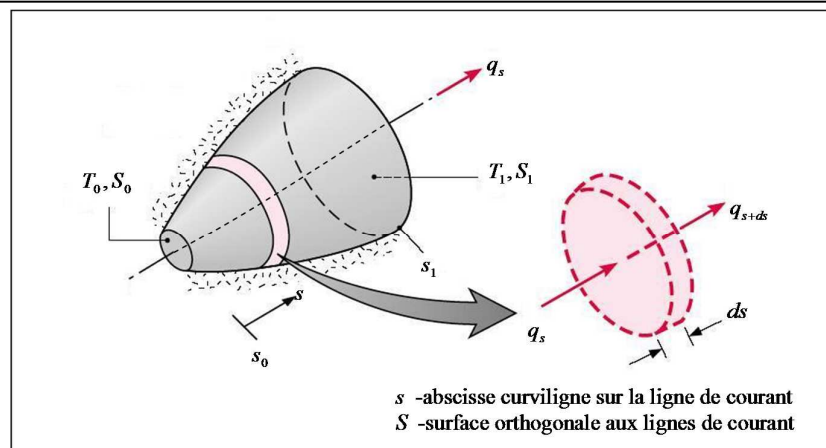
$$q_x = A\varphi = \lambda \frac{A}{L}(T_{S1} - T_{S2})$$

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Steady-state conduction without internal sources

• Fourier law in a section

$$q = S(s) \varphi = -S(s) \lambda(s) \frac{dT}{ds} \Big|_s$$

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– Separation of variables

$$q \frac{ds}{S(s) \lambda(s)} = -dT$$

$$q \int_{s_0}^s \frac{ds}{S(s) \lambda(s)} = - \int_{T_0}^T dT \quad \Leftrightarrow \quad qR = T_0 - T_1$$

$$R = \int_{s_0}^s \frac{ds}{S(s) \lambda(s)}$$

$$R_{cd} \equiv \frac{T_0 - T_1}{q}$$

$$R_{cd} = \frac{1}{\lambda} \frac{L}{A}$$

Plane wall

$$R_{cd} = \frac{1}{\lambda} \frac{\ln(r_2 / r_1)}{2\pi L}$$

cylinder

$$R_{cd} = \frac{1}{\lambda} \frac{1}{4\pi} \left(\frac{1}{r_1} - \frac{1}{r_2} \right)$$

sphere

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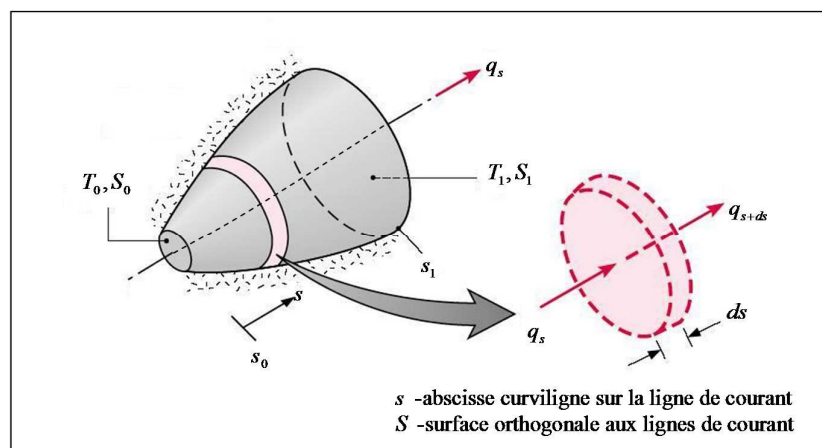
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• Stream surface (tube)



$$q_s - q_{s+ds} + p S ds = 0 \quad \text{Steady-state conduction}$$

$$dq = p S ds$$

$$\int_{q_0}^q dq = \int_{s_0}^s p S ds$$

$$q(s) = \int_{s_0}^s p S ds + q_0$$

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- stream surface (tube)

$$q(s) = -\lambda S \frac{dT}{ds}$$

$$q(s) = \int_{s_0}^s p S ds + q_0$$

$$\lambda S \frac{dT}{ds} = - \int_{s_0}^s p S ds - q_0$$

$$dT = \frac{1}{\lambda S} \left[- \int_{s_0}^s p S ds' \right] ds - q_0 \frac{ds}{\lambda S}$$

$$T_1 = \int_{s_0}^{s_1} \frac{1}{\lambda S} \left[\int_{s_0}^s - p S ds' \right] ds - q_0 \int_{s_0}^{s_1} \frac{ds}{\lambda S} + T_0$$

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$$dr = \frac{ds}{\lambda S}$$

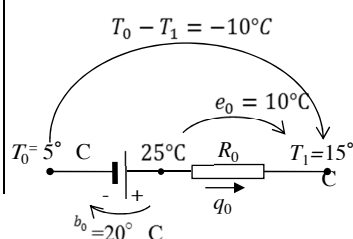
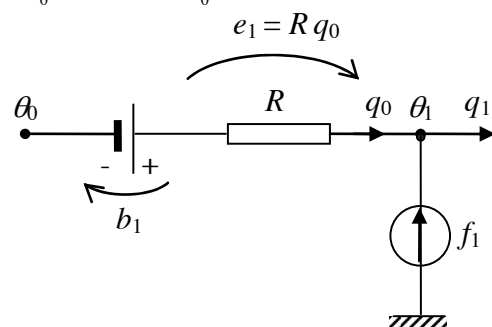
$$R = \int_{s_0}^{s_1} \frac{ds}{\lambda S}$$

$$T_1 = -r_1 \int_{s_0}^{s_1} p S ds + \int_{s_0}^{s_1} p S r ds - R q_0 + T_0$$

$$f_1 = \int_{s_0}^{s_1} p S ds$$

$$b_1 = -r_1 \int_{s_0}^{s_1} p S ds + \int_{s_0}^{s_1} p S r ds$$

$$\begin{cases} q_1 = q_0 + f_1 \\ \theta_1 = b_1 + e_1 + \theta_0 \end{cases}$$



$$e_0 = T_0 + b_0 - T_1$$

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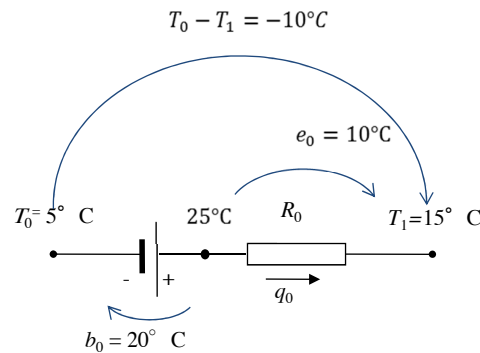
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$$e_0 = T_0 + b_0 - T_1$$

- Convection resistance

$$R_{cv} \equiv \frac{T_s - T_\infty}{q} = \frac{1}{h_{cv} A}$$

- Radiative resistance

$$R_r \equiv \frac{T_s - T_{me}}{q} = \frac{1}{h_r A}$$

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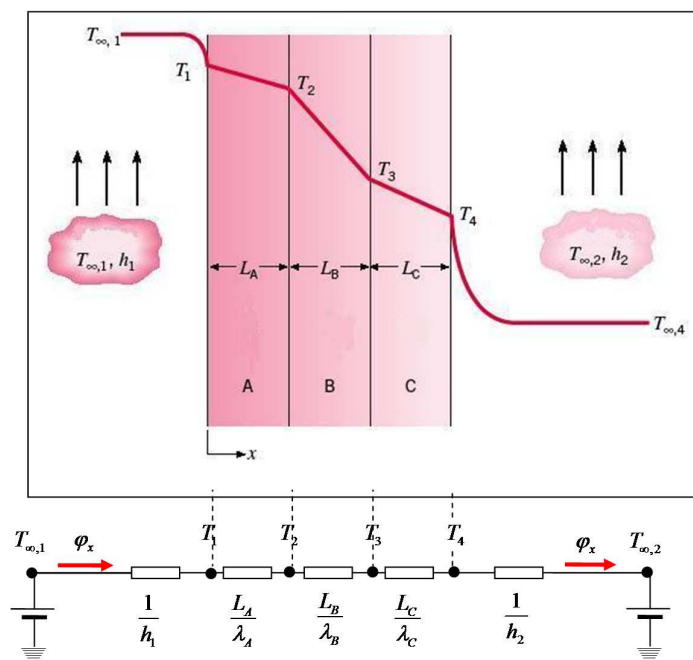
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- Multilayer 1-D wall



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• Mur multicouche

$$R_{tot} = \frac{1}{h_1 A} + \frac{L_A}{\lambda_A A} + \frac{L_B}{\lambda_B A} + \frac{L_C}{\lambda_C A} + \frac{1}{h_2 A}$$

$$q_x \equiv U \cdot A \cdot \Delta T$$

$$U = \frac{1}{\frac{1}{h_1} + \frac{L_A}{\lambda_A} + \frac{L_B}{\lambda_B} + \frac{L_C}{\lambda_C} + \frac{1}{h_2}}$$

$$q_x = \frac{T_{\infty 1} - T_{\infty 2}}{R_{tot}} = UA(T_{\infty 1} - T_{\infty 2})$$

$$\varphi_x = U(T_{\infty 1} - T_{\infty 2}) = \frac{T_{\infty 1} - T_{\infty 2}}{\frac{1}{h_1} + \frac{L_A}{\lambda_A} + \frac{L_B}{\lambda_B} + \frac{L_C}{\lambda_C} + \frac{1}{h_2}}$$

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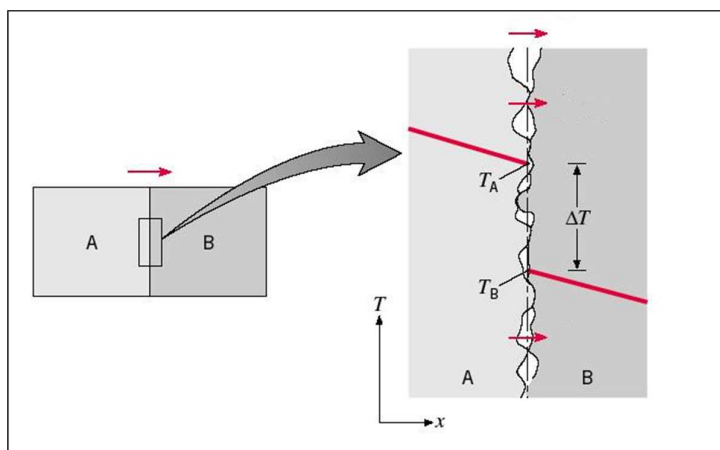
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• Contact resistance



$$R_{ct} \equiv \frac{T_A - T_B}{q_x}$$

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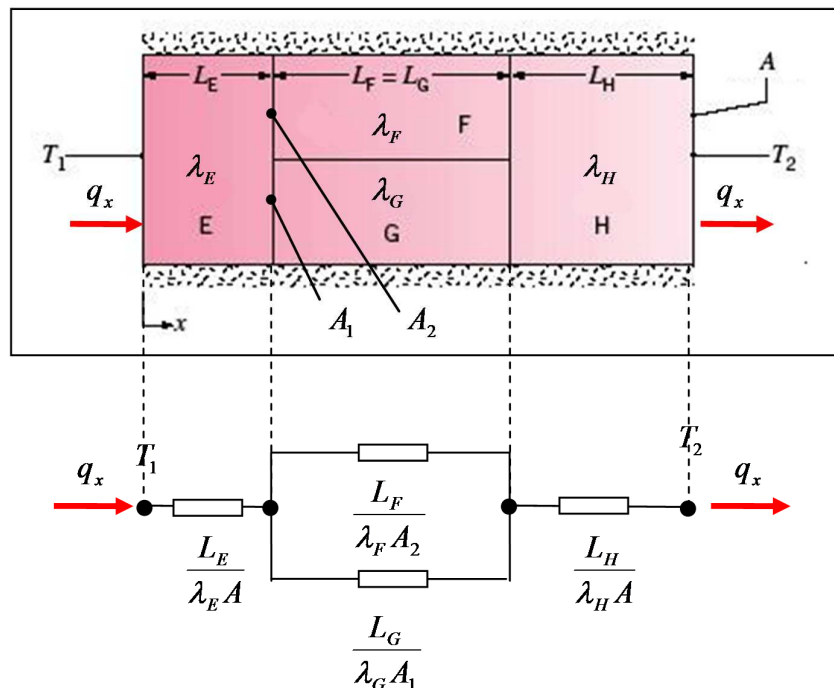
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• Example circuit



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Type de condition	Equation	Variation de la température	Source équivalente
1. Température imposée sur la surface (condition de Dirichlet)	$T(0, t) = T_s$		
2. Densité de flux imposée sur la surface (condition de Neumann)			
a. Densité de flux imposée	$-\lambda \frac{\partial T}{\partial x} \Big _{x=0} = \varphi_s$		
b. Surface adiabatique ou surface de symétrie	$\frac{\partial T}{\partial x} \Big _{x=0} = \varphi_s = 0$		
3. Surface avec convection (condition de Fourier)	$-\lambda \frac{\partial T}{\partial x} \Big _{x=0} = h(T_\infty - T(0, t))$ $= h(T_\infty - T_s)$		

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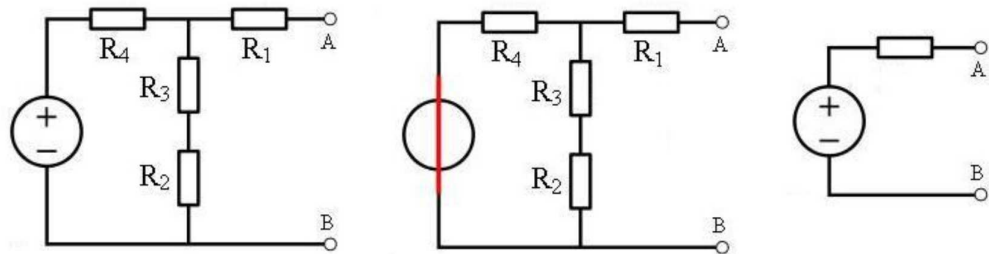
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Thévenin theorem

- tension $\leftarrow$ potential diff.
- resistance $\leftarrow$ passivized sources



$$q_{tot} = \frac{T}{R_1 + \left(\frac{1}{R_2} + \frac{1}{R_3} \right)^{-1}}$$

$$T_T = \frac{R_2 + R_3}{(R_2 + R_3) + R_4} T_1$$

$$R_T = R_1 + \left(\frac{1}{R_2 + R_3} + \frac{1}{R_4} \right)^{-1}$$

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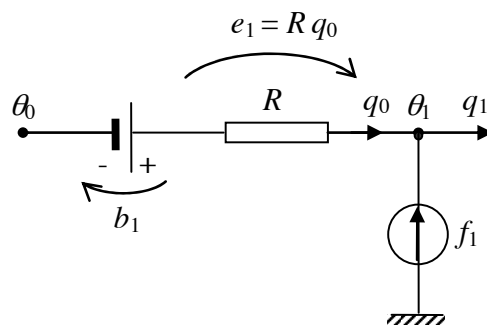
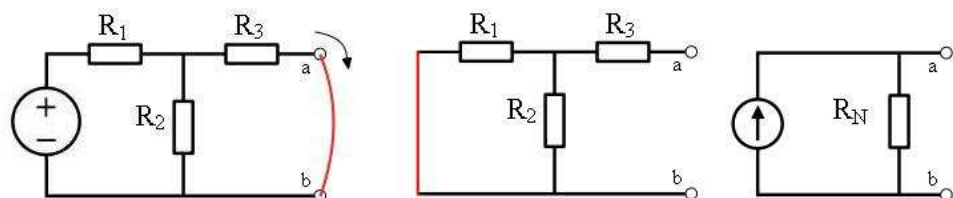
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Norton theorem

- current $\leftarrow$ short circuit terminals
- resistance $\leftarrow$ passivized sources



$$q_N = \frac{R_2}{R_2 + R_3} q_{tot}$$

$$R_N = R_3 + \left(\frac{1}{R_1} + \frac{1}{R_2} \right)^{-1}$$

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• Heat equation

$$\rho c \frac{\partial T}{\partial t} = \lambda \operatorname{div}(\mathbf{grad} T) + p$$

• Homogenous and isotropic material

$$\lambda = \text{const.}$$

steady-state

$$\frac{\partial T}{\partial t} = 0$$

Poisson equation:

$$\lambda \operatorname{div}(\mathbf{grad} T) + p = 0$$

$$\lambda \Delta T + p = 0$$

Conduction

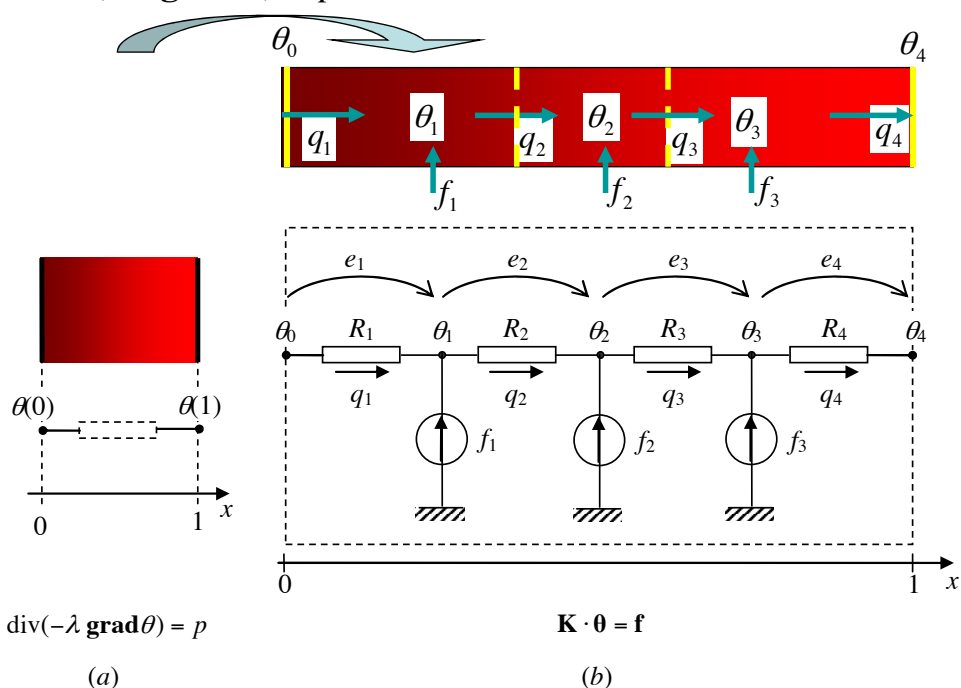
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• Poisson equation

$$\operatorname{div}(-\lambda \mathbf{grad} \theta) = p$$



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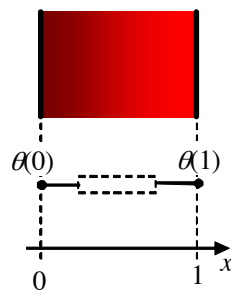
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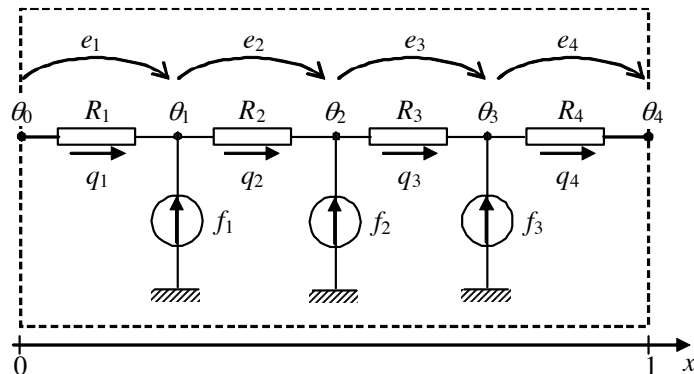
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• Direct problem



$$\text{div}(\lambda \text{grad} \theta) = -p$$

(a)



$$\mathbf{A}^T \mathbf{C} \mathbf{A} \cdot \boldsymbol{\theta} = -\mathbf{f}$$

(b)

- **given:** circuit and boundary conditions
- **find:** node temperatures and heat flow rates on branches

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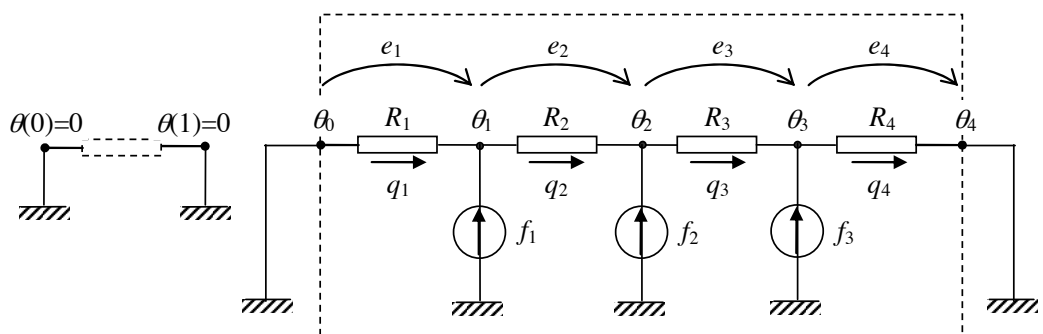
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• Dirichlet boundary conditions



– temperature differences for each resistance

$$\begin{cases} e_1 = -\theta_1 \\ e_2 = \theta_1 - \theta_2 \\ e_3 = \theta_2 - \theta_3 \\ e_4 = \theta_3 \end{cases} \quad \begin{bmatrix} e_1 \\ e_2 \\ e_3 \\ e_4 \end{bmatrix} = \begin{bmatrix} -1 & 0 & 0 \\ 1 & -1 & 0 \\ 0 & 1 & -1 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} \theta_1 \\ \theta_2 \\ \theta_3 \end{bmatrix} \quad \mathbf{e} = -\mathbf{A} \cdot \boldsymbol{\theta}$$

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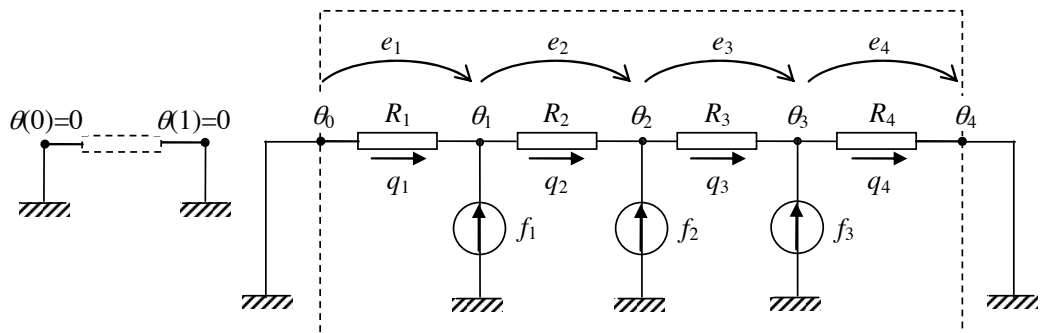
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• Dirichlet boundary conditions



– heat flux : constitutive law (Fourier)

$$\begin{cases} q_1 = G_1 e_1 \\ q_2 = G_2 e_2 \\ q_3 = G_3 e_3 \\ q_4 = G_4 e_4 \end{cases} \quad \begin{bmatrix} q_1 \\ q_2 \\ q_3 \\ q_4 \end{bmatrix} = \begin{bmatrix} G_1 & 0 & 0 & 0 \\ 0 & G_2 & 0 & 0 \\ 0 & 0 & G_3 & 0 \\ 0 & 0 & 0 & G_4 \end{bmatrix} \begin{bmatrix} e_1 \\ e_2 \\ e_3 \\ e_4 \end{bmatrix} \quad \mathbf{q} = \mathbf{G} \cdot \mathbf{e}$$

$$G_i = 1 / R_i, i = 1, \dots, 4$$

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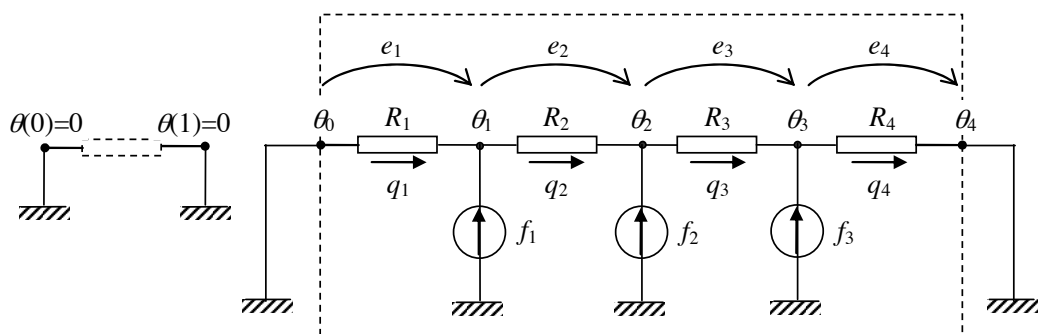
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• Dirichlet boundary conditions



– heat balance in nodes

$$\begin{cases} q_1 - q_2 = -f_1 \\ q_2 - q_3 = -f_2 \\ q_3 - q_4 = -f_3 \end{cases} \quad \begin{bmatrix} f_1 \\ f_2 \\ f_3 \end{bmatrix} = \begin{bmatrix} 1 & -1 & 0 & 0 \\ 0 & 1 & -1 & 0 \\ 0 & 0 & 1 & -1 \end{bmatrix} \begin{bmatrix} q_1 \\ q_2 \\ q_3 \\ q_4 \end{bmatrix} \quad \mathbf{f} = -\mathbf{A}^T \cdot \mathbf{q}$$

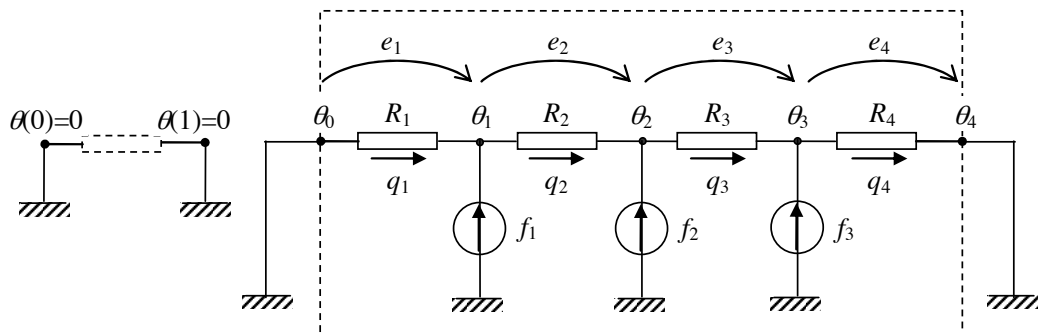
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• Dirichlet boundary conditions



– solution

$$\mathbf{K} \cdot \boldsymbol{\theta} = \mathbf{f} \quad \mathbf{K} = \mathbf{A}^T \mathbf{G} \mathbf{A} = \begin{bmatrix} G_1 + G_2 & -G_2 & 0 \\ -G_2 & G_2 + G_3 & -G_3 \\ 0 & -G_3 & G_3 + G_4 \end{bmatrix}$$

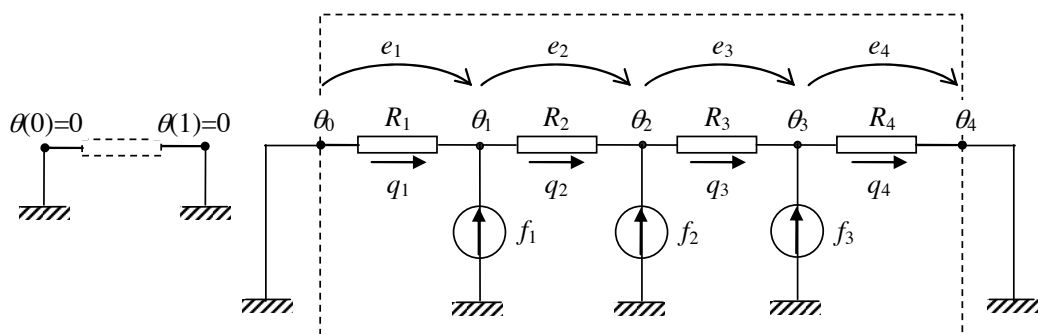
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• Dirichlet boundary conditions



$$\mathbf{f} = \mathbf{K} \cdot \boldsymbol{\theta} \quad \longrightarrow \quad \boldsymbol{\theta} = \mathbf{K}^{-1} \cdot \mathbf{f}$$

$$\mathbf{G} = \mathbf{I}$$

$$\mathbf{K} \equiv \mathbf{A}^T \mathbf{A} = \begin{bmatrix} 1 & 0 & 0 \\ -1 & 1 & 0 \\ 0 & -1 & 1 \\ 0 & 0 & -1 \end{bmatrix} \cdot \begin{bmatrix} 1 & -1 & 0 & 0 \\ 0 & 1 & -1 & 0 \\ 0 & 0 & 1 & -1 \end{bmatrix} = \begin{bmatrix} 2 & -1 & 0 \\ -1 & 2 & -1 \\ 0 & -1 & 2 \end{bmatrix}$$

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• Dirichlet boundary conditions

$$\mathbf{K} \equiv \mathbf{A}^T \mathbf{A} = \begin{bmatrix} 2 & -1 & 0 \\ -1 & 2 & -1 \\ 0 & -1 & 2 \end{bmatrix}$$

• Properties of K matrix

- symmetric $\mathbf{K} = \mathbf{K}^T$
- tri-diagonal
- sparse
- constant diagonal

• Toeplitz matrix

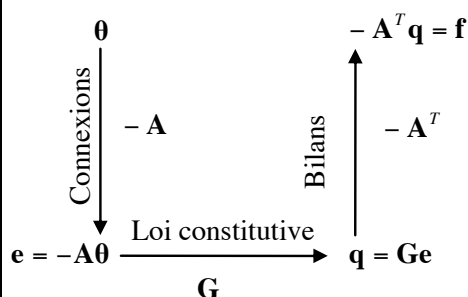
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• Framework for analysis of thermal circuits



θ températures dans les nœuds
 e chute de température sur les résistances
 q flux dans les résistances
 f flux externes

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• Equivalence between continuous and discrete

	$T \quad \text{Connexion} \quad \nabla T \quad \text{grad} \quad \text{Loi constitutive} \quad -\lambda \quad \varphi \quad \text{Bilan} \quad \text{div} \quad f = p - \rho c \frac{\partial T}{\partial t}$	$\theta \quad \text{Connexions} \quad -A \quad \text{Loi constitutive} \quad G \quad q = Ge \quad \text{Bilans} \quad -A^T \quad A^T q = f$
	Continu	Discret
Equation de la chaleur	$\text{div}(\lambda \text{ grad } \theta) = -p$	$A^T G A \theta = f$
Conditions aux limites	en $x = 0 : \theta(0)$ ou $\varphi(0)$ en $x = 1 : \theta(1)$ ou $\varphi(1)$	incluses dans la matrice $A^T G A$
Température	θ	θ
Flux	φ	q
Connexion	$-\text{grad } \theta$	$-A\theta$
Bilan d'énergie	$-\text{div } \varphi$	$-A^T q$
Constante de matériau	λ	G
Sources	p	$-f$

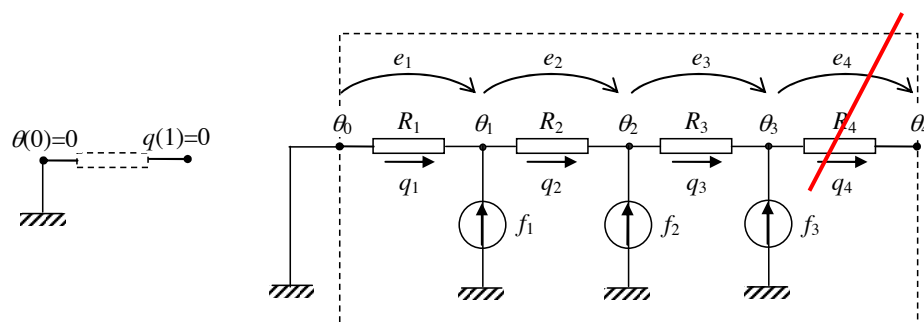
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• Dirichlet plus Neumann



– Connectivity matrix: circuit topology

nœud θ_1 θ_2 θ_3 arc

$$A = \begin{bmatrix} 1 & 0 & 0 \\ -1 & 1 & 0 \\ 0 & -1 & 1 \end{bmatrix} \begin{matrix} q_1 \\ q_2 \\ q_3 \\ \underline{q_4} \end{matrix}$$

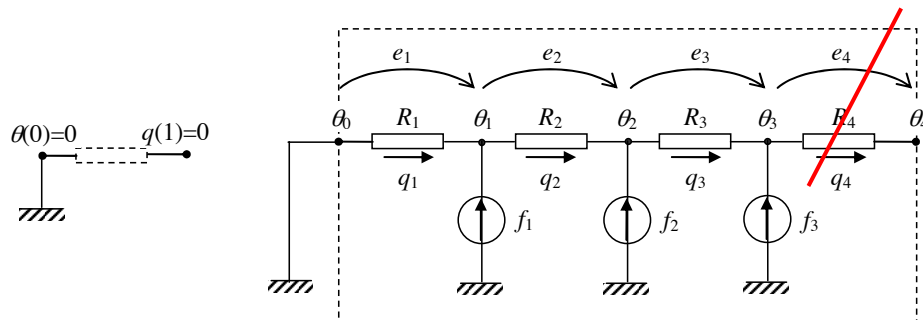
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• Dirichlet plus Neumann



– flux: constitutive law

$$\begin{cases} q_1 = G_1 e_1 \\ q_2 = G_2 e_2 \\ q_3 = G_3 e_3 \end{cases} \quad \begin{bmatrix} q_1 \\ q_2 \\ q_3 \end{bmatrix} = \begin{bmatrix} G_1 & 0 & 0 \\ 0 & G_2 & 0 \\ 0 & 0 & G_3 \end{bmatrix} \begin{bmatrix} e_1 \\ e_2 \\ e_3 \end{bmatrix} \quad \mathbf{q} = \mathbf{G} \cdot \mathbf{e}$$

$$q_4 = G_4 e_4 = 0 \quad (\text{connu})$$

$$G_i = 1 / R_i, i = 1, \dots, 4$$

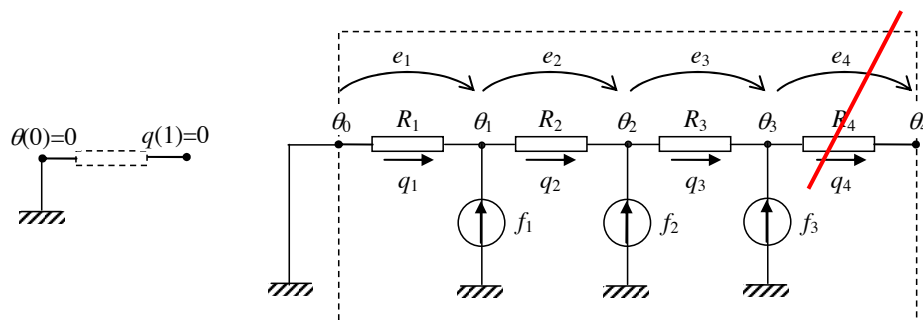
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• Dirichlet plus Neumann



– heat balance

$$\begin{cases} q_1 - q_2 = -f_1 \\ q_2 - q_3 = -f_2 \\ q_3 - q_4 = -f_3 \end{cases} \quad \begin{bmatrix} f_1 \\ f_2 \\ f_3 \end{bmatrix} = - \begin{bmatrix} 1 & -1 & 0 \\ 0 & 1 & -1 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} q_1 \\ q_2 \\ q_3 \end{bmatrix} \quad \mathbf{f} = -\mathbf{A}^T \cdot \mathbf{q}$$

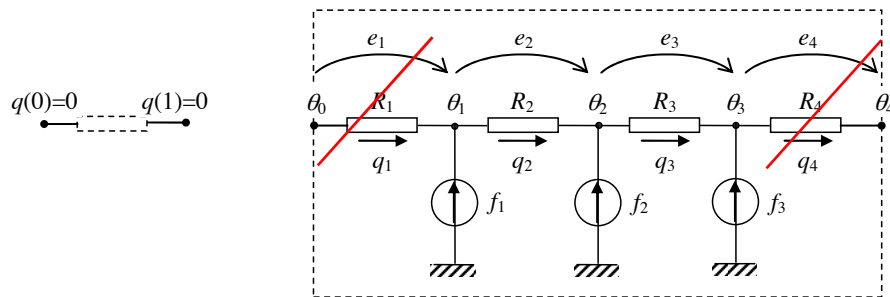
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• Neumann conditions at both boundaries



– connectivity matrix

$$A = \begin{bmatrix} -1 & 1 & 0 \\ 0 & -1 & 1 \end{bmatrix} \begin{matrix} R_1 \\ R_2 \\ R_3 \\ R_4 \end{matrix}$$

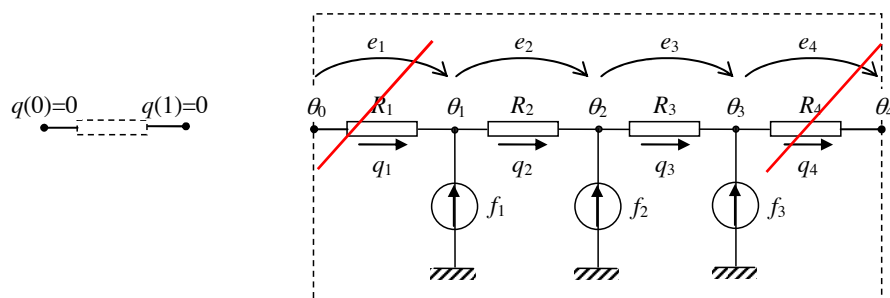
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• Neumann conditions at both boundaries



– constitutive law

$$\begin{cases} q_1 = 0 \\ q_2 = G_2 e_2 \\ q_3 = G_3 e_3 \\ q_4 = 0 \end{cases} \quad \begin{bmatrix} q_2 \\ q_3 \end{bmatrix} = \begin{bmatrix} G_2 & 0 \\ 0 & G_3 \end{bmatrix} \begin{bmatrix} e_2 \\ e_3 \end{bmatrix} \quad \mathbf{q} = \mathbf{G} \cdot \mathbf{e}$$

$$G_i = 1/R_i, i = 1, \dots, 4$$

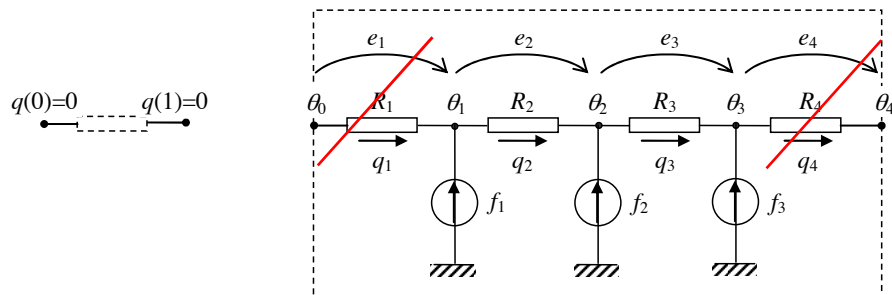
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• Neumann conditions at both boundaries



– heat balance

$$\begin{cases} -q_2 = -f_1 \\ q_2 - q_3 = -f_2 \\ q_3 = -f_3 \end{cases} \quad \begin{bmatrix} f_1 \\ f_2 \\ f_3 \end{bmatrix} = - \begin{bmatrix} -1 & 0 \\ 1 & -1 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} q_2 \\ q_3 \end{bmatrix} \quad \mathbf{f} = -\mathbf{A}^T \cdot \mathbf{q}$$

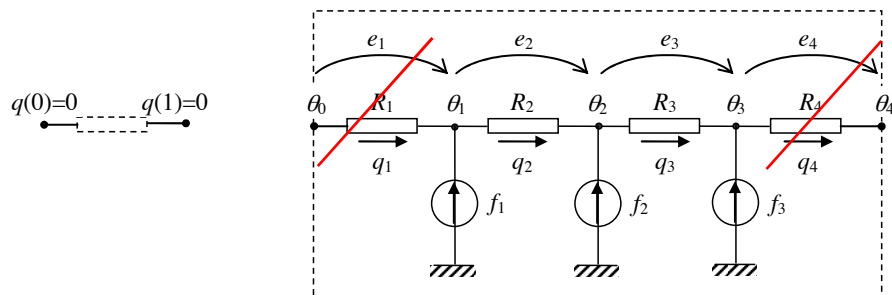
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• Neumann conditions at both boundaries



– Circuit equation

$$\mathbf{f} = \mathbf{B} \cdot \boldsymbol{\theta} \quad \mathbf{C} = \mathbf{I} \quad \mathbf{B} = \mathbf{A}^T \mathbf{C} \mathbf{A} = \begin{bmatrix} 1 & -1 & 0 \\ -1 & 2 & -1 \\ 0 & -1 & 1 \end{bmatrix}$$

singulière ou non inversable

To solve: temperature reference is needed

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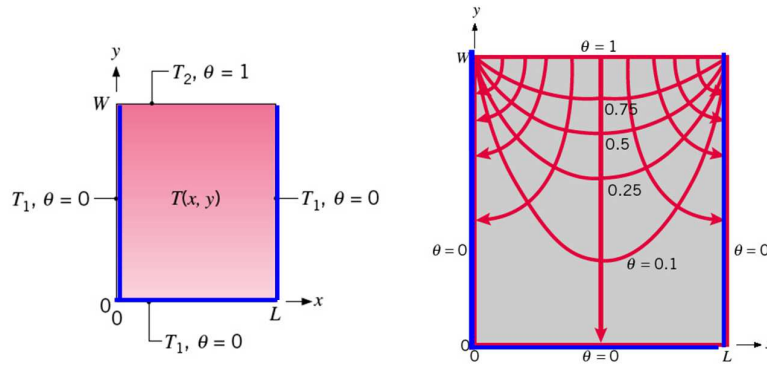
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- 2D, steady-state, without internal sources



$$\rho c \frac{\partial T}{\partial t} = \text{div}(\lambda \mathbf{grad} T) + p \quad \longrightarrow \quad \frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} = 0$$

Eq. Laplace

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- **Finite differences**
 - Discretization of derivative
- **Finite volume**
 - mean values of conservative variables in a volume, not at the nodes (by integration)
 - non-invasive boundary conditions
 - structured or unstructured meshing
- **Finite element**
 - use a function on a domain
 - unstructured meshing

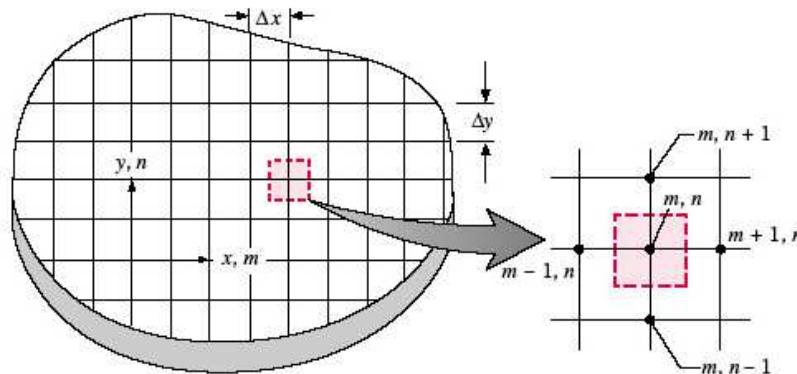
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Energy Building Simulation
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- Finite volume
 - volume discretization



- energy balance for every control volume

$$\dot{E}_e - \dot{E}_s + \dot{E}_g = 0$$

Conduction

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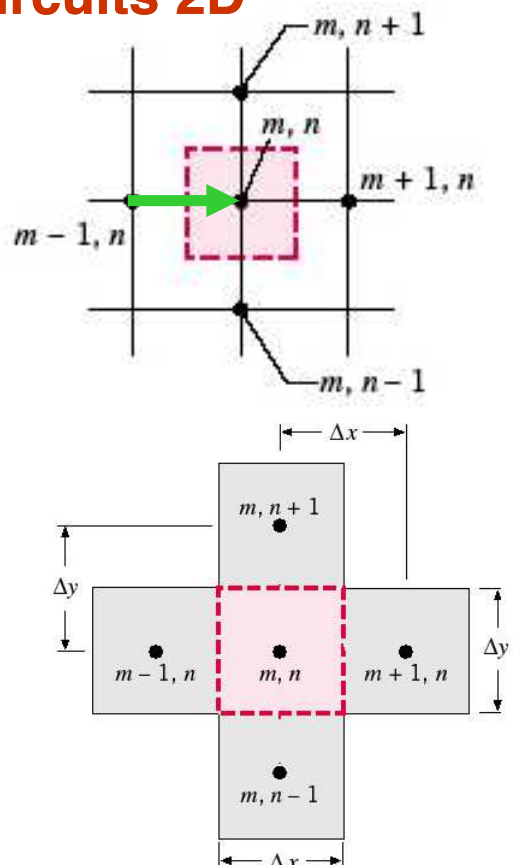
Energy Building Simulation
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- flux

$$q_x = -\lambda A \frac{dT}{dx}$$

$$q_{(m-1,n) \rightarrow (m,n)} = \underbrace{\lambda(\Delta y \cdot 1)}_A \underbrace{\frac{\theta_{m-1,n} - \theta_{m,n}}{\Delta x}}_{-\text{grad} \theta}$$

$$R_{i,j} = \frac{1}{\lambda} \frac{L}{A} = \frac{1}{\lambda} \frac{\Delta x}{\Delta y \cdot 1}$$



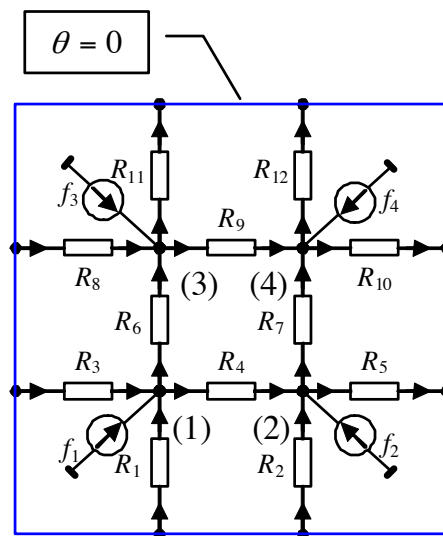
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• Dirichlet boundary conditions



$$\mathbf{A} = \begin{matrix} \text{nœud} & \theta_1 & \theta_2 & \theta_3 & \theta_4 & \text{arc} \\ \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ -1 & 1 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ -1 & 0 & 1 & 0 \\ 0 & -1 & 0 & 1 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & -1 & 1 \\ 0 & 0 & 0 & -1 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{bmatrix} & R_1 \\ R_2 \\ R_3 \\ R_4 \\ R_5 \\ R_6 \\ R_7 \\ R_8 \\ R_9 \\ R_{10} \\ R_{11} \\ R_{12} \end{matrix}$$

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• Dirichlet boundary conditions

– Circuit equation

$$\mathbf{f} = \mathbf{K} \cdot \boldsymbol{\theta}$$

$$f_i = p(\Delta x \cdot \Delta y \cdot 1) \quad p [\text{W/m}^3]$$

$$\mathbf{K} = \mathbf{A}^T \mathbf{G} \mathbf{A}$$

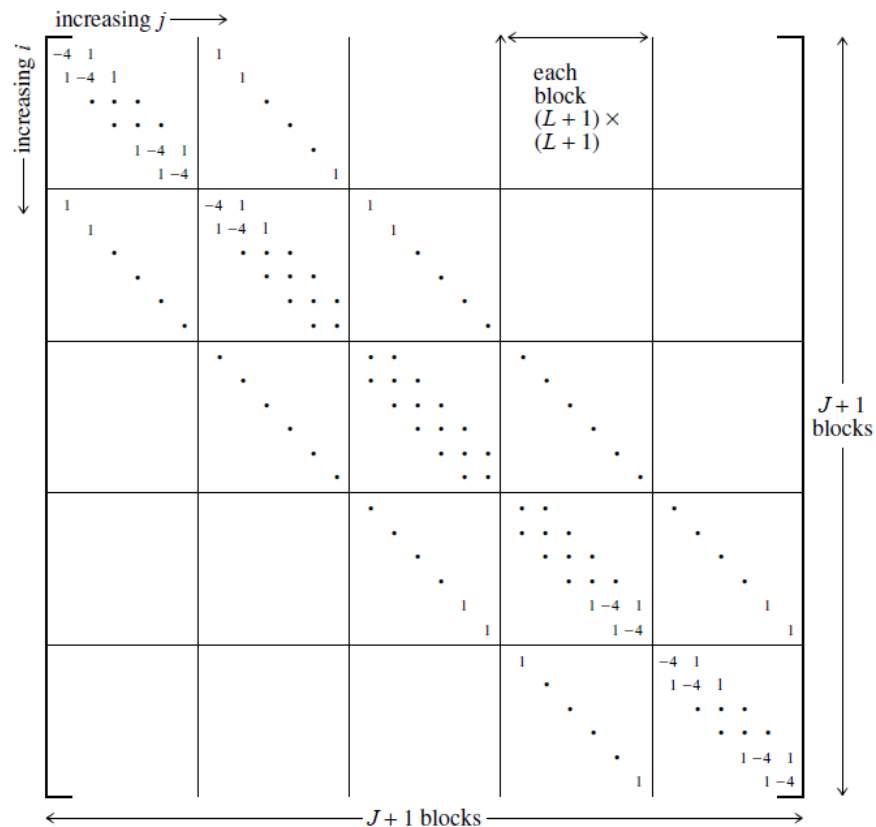
$$\mathbf{K} = \mathbf{A}^T \mathbf{A} = \begin{bmatrix} 4 & -1 & -1 & 0 \\ -1 & 4 & 0 & -1 \\ -1 & 0 & 4 & -1 \\ 0 & -1 & -1 & 4 \end{bmatrix}$$

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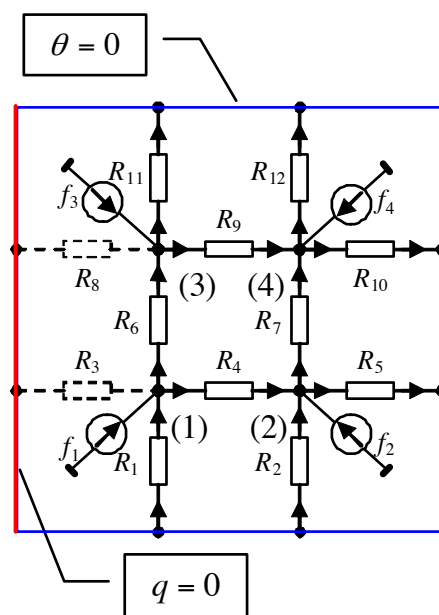
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• Dirichlet and Neumann boundary conditions



$$A = \begin{matrix} \text{noeud} & \theta_1 & \theta_2 & \theta_3 & \theta_4 & \text{arc} \\ \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ \boxed{0} & 0 & 0 & 0 \\ -1 & 1 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ -1 & 0 & 1 & 0 \\ 0 & -1 & 0 & 1 \\ 0 & 0 & \boxed{0} & 0 \\ 0 & 0 & -1 & 1 \\ 0 & 0 & 0 & -1 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{bmatrix} & R_1 \\ R_2 \\ R_3 \\ R_4 \\ R_5 \\ R_6 \\ R_7 \\ R_8 \\ R_9 \\ R_{10} \\ R_{11} \\ R_{12} \end{matrix}$$

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- Dirichlet and Neumann boundary conditions
 - circuit equations

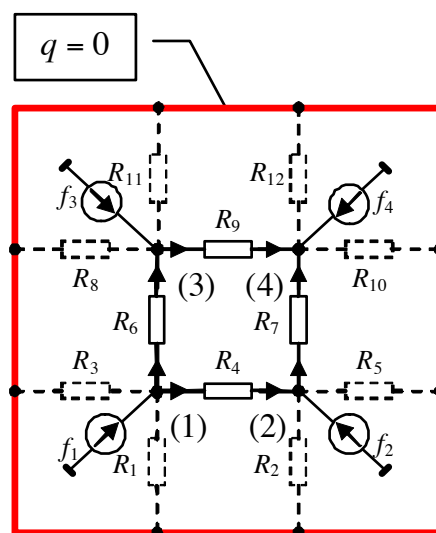
$$\mathbf{K} = \mathbf{A}^T \mathbf{A} = \begin{bmatrix} 3 & -1 & -1 & 0 \\ -1 & 4 & 0 & -1 \\ -1 & 0 & 3 & -1 \\ 0 & -1 & -1 & 4 \end{bmatrix}$$

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- Neumann boundary conditions



nœud	θ_1	θ_2	θ_3	θ_4	arc
	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	R_1
	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	R_2
	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	R_3
	$\begin{bmatrix} -1 \\ 0 \\ -1 \end{bmatrix}$	$\begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix}$	$\begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix}$	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	R_4
	$\begin{bmatrix} 0 \\ -1 \\ 0 \end{bmatrix}$	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	$\begin{bmatrix} 0 \\ 1 \\ 0 \end{bmatrix}$	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	R_5
	$\begin{bmatrix} -1 \\ 0 \\ 0 \end{bmatrix}$	$\begin{bmatrix} 0 \\ -1 \\ 0 \end{bmatrix}$	$\begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix}$	$\begin{bmatrix} 0 \\ 0 \\ 1 \end{bmatrix}$	R_6
	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	$\begin{bmatrix} -1 \\ 0 \\ 0 \end{bmatrix}$	$\begin{bmatrix} 0 \\ 1 \\ 0 \end{bmatrix}$	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	R_7
	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	R_8
	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	$\begin{bmatrix} -1 \\ 0 \\ 0 \end{bmatrix}$	$\begin{bmatrix} 1 \\ 0 \\ 0 \end{bmatrix}$	R_9
	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	R_{10}
	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	R_{11}
	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	$\begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$	R_{12}

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• Neumann boundary conditions

$$\mathbf{K} = \mathbf{A}^T \mathbf{A} = \begin{bmatrix} 2 & -1 & -1 & 0 \\ -1 & 2 & 0 & -1 \\ -1 & 0 & 2 & -1 \\ 0 & -1 & -1 & 2 \end{bmatrix}$$

singular matrix

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$$\begin{bmatrix} 4 & -1 & -1 & 0 \\ -1 & 4 & 0 & -1 \\ -1 & 0 & 4 & -1 \\ 0 & -1 & -1 & 4 \end{bmatrix}$$

Dirichlet

$$\begin{bmatrix} 3 & -1 & -1 & 0 \\ -1 & 4 & 0 & -1 \\ -1 & 0 & 3 & -1 \\ 0 & -1 & -1 & 4 \end{bmatrix}$$

Dirichlet & Neumann

$$\begin{bmatrix} 2 & -1 & -1 & 0 \\ -1 & 2 & 0 & -1 \\ -1 & 0 & 2 & -1 \\ 0 & -1 & -1 & 2 \end{bmatrix}$$

Neumann

$$\frac{\partial^2 \theta}{\partial x^2} \bigg|_{m,n} = \frac{\frac{\partial \theta}{\partial x} \bigg|_{m+1/2,n} - \frac{\partial \theta}{\partial x} \bigg|_{m-1/2,n}}{\Delta x}$$

$$= \frac{\theta_{m-1,n} - 2\theta_{m,n} + \theta_{m+1,n}}{(\Delta x)^2}$$

$$\frac{\partial^2 \theta}{\partial y^2} \bigg|_{m,n} = \frac{\theta_{m,n-1} - 2\theta_{m,n} + \theta_{m,n+1}}{(\Delta y)^2}$$

$$-(\Delta x)^2 \left(\frac{\partial^2 \theta}{\partial x^2} + \frac{\partial^2 \theta}{\partial y^2} \right) = 4\theta_{m,n} - \theta_{m,n-1} - \theta_{m,n+1} - \theta_{m-1,n} - \theta_{m+1,n}$$

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• heat equation

$$\rho c \frac{\partial T}{\partial t} = -\text{div}(-\lambda \mathbf{grad} T) + p$$

• 1D

$$\rho c \frac{\partial T}{\partial t} = \lambda \frac{d^2 T}{dx^2} + p$$

• finite differences 1D

$$\rho \cdot c \dot{\theta}_m = \lambda \frac{1}{\Delta x^2} [(\theta_{m-1} - \theta_m) + (\theta_{m+1} - \theta_m)] + p$$

$$\rho \cdot A \Delta x \cdot c \dot{\theta}_m = \lambda \frac{A}{\Delta x} (\theta_{m-1} - \theta_m) + \lambda \frac{A}{\Delta x} (\theta_{m+1} - \theta_m) + \dot{q} \cdot A \Delta x$$

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• thermal circuit 1D

$$\rho c \frac{\partial T}{\partial t} = -\text{div}(-\lambda \mathbf{grad} T) + p$$

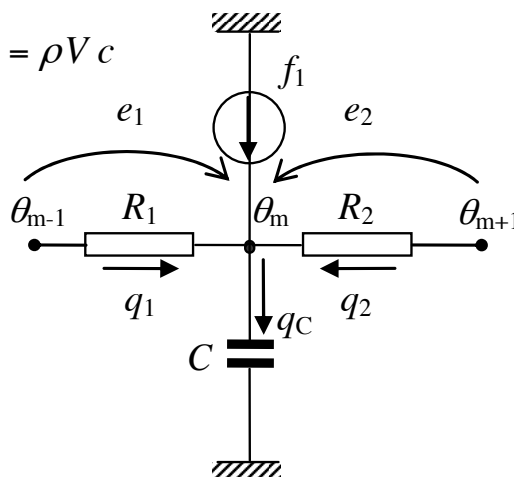
$$\rho \cdot A \Delta x \cdot c \dot{\theta}_m = \lambda \frac{A}{\Delta x} (\theta_{m-1} - \theta_m) + \lambda \frac{A}{\Delta x} (\theta_{m+1} - \theta_m) + \dot{q} \cdot A \Delta x$$

$$C \equiv \rho A \Delta x c = \rho V c$$

$$R_1 \equiv \lambda \frac{A}{\Delta x}$$

$$R_2 \equiv \lambda \frac{A}{\Delta x}$$

$$f_1 \equiv \dot{q} A \Delta x$$



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- **het equation**

- **steady-state**

$$0 = -\text{div}(-\lambda \text{ grad} T) + p$$

$$0 = -\mathbf{A}^T \mathbf{G} \mathbf{A} \boldsymbol{\theta} + \mathbf{f}_T ; \quad \mathbf{f}_T \equiv \mathbf{A}^T \mathbf{G} \mathbf{b} + \mathbf{f}$$

- **dynamic**

$$\rho c \frac{\partial T}{\partial t} = -\text{div}(-\lambda \text{ grad} T) + p$$

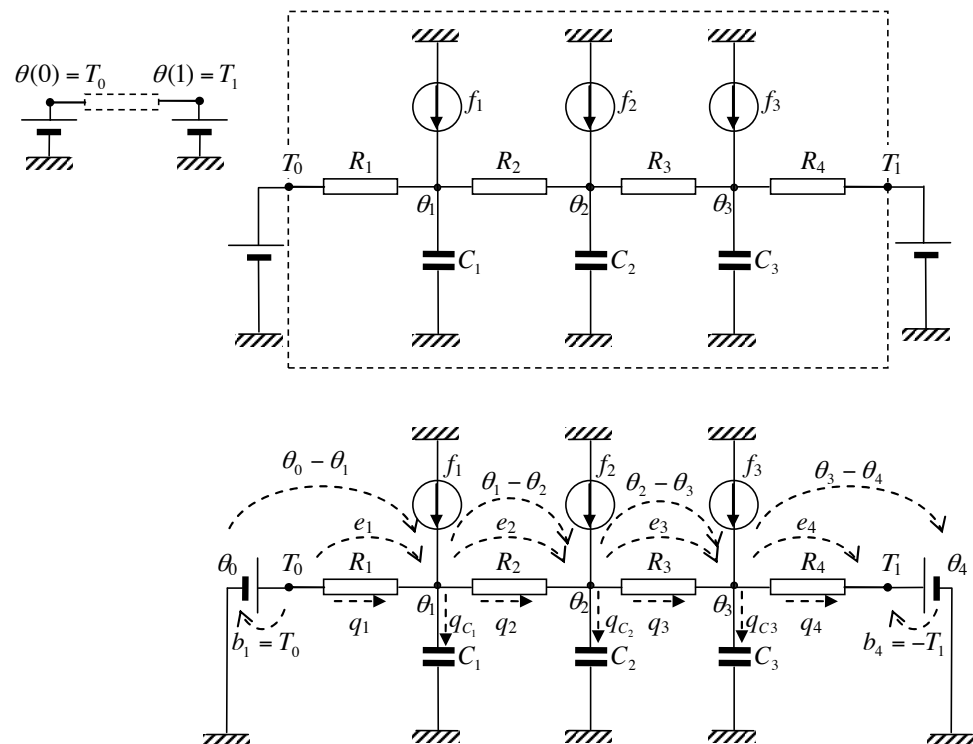
$$\mathbf{C} \dot{\boldsymbol{\theta}} = -\mathbf{A}^T \mathbf{G} \mathbf{A} \cdot \boldsymbol{\theta} + \mathbf{f}_T ; \quad \mathbf{f}_T \equiv \mathbf{A}^T \mathbf{G} \mathbf{b} + \mathbf{f}$$

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$$\begin{cases} e_1 - b_1 = \theta_0 - \theta_1 \\ e_2 - b_2 = \theta_1 - \theta_2 \\ e_3 - b_3 = \theta_2 - \theta_3 \\ e_4 - b_4 = \theta_3 - \theta_4 \end{cases} \quad \begin{bmatrix} e_1 \\ e_2 \\ e_3 \\ e_4 \end{bmatrix} - \begin{bmatrix} b_1 \\ 0 \\ 0 \\ b_4 \end{bmatrix} = \begin{bmatrix} -1 & 0 & 0 \\ 1 & -1 & 0 \\ 0 & 1 & -1 \\ 0 & 0 & 1 \end{bmatrix} \begin{bmatrix} \theta_1 \\ \theta_2 \\ \theta_3 \end{bmatrix} \quad \mathbf{e} - \mathbf{b} = -\mathbf{A}\boldsymbol{\theta}$$

$$\begin{cases} q_1 = G_1 e_1 \\ q_2 = G_2 e_2 \\ q_3 = G_3 e_3 \\ q_4 = G_4 e_4 \end{cases} \quad \begin{bmatrix} q_1 \\ q_2 \\ q_3 \\ q_4 \end{bmatrix} = \begin{bmatrix} G_1 & 0 & 0 & 0 \\ 0 & G_2 & 0 & 0 \\ 0 & 0 & G_3 & 0 \\ 0 & 0 & 0 & G_4 \end{bmatrix} \begin{bmatrix} e_1 \\ e_2 \\ e_3 \\ e_4 \end{bmatrix} \quad \mathbf{q} = \mathbf{G} \cdot \mathbf{e}$$

$$\begin{cases} C_1 \dot{\theta}_1 = q_1 - q_2 + f_1 \\ C_2 \dot{\theta}_2 = q_2 - q_3 + f_2 \\ C_3 \dot{\theta}_3 = q_3 - q_4 + f_3 \end{cases} \quad \begin{bmatrix} C_1 & 0 & 0 \\ 0 & C_2 & 0 \\ 0 & 0 & C_3 \end{bmatrix} \begin{bmatrix} \dot{\theta}_1 \\ \dot{\theta}_2 \\ \dot{\theta}_3 \end{bmatrix} = \begin{bmatrix} 1 & -1 & 0 & 0 \\ 0 & 1 & -1 & 0 \\ 0 & 0 & 1 & -1 \end{bmatrix} \begin{bmatrix} q_1 \\ q_2 \\ q_3 \\ q_4 \end{bmatrix} + \begin{bmatrix} f_1 \\ f_2 \\ f_3 \end{bmatrix} \quad \mathbf{C}\dot{\boldsymbol{\theta}} = \mathbf{A}^T \mathbf{q} + \mathbf{f}$$

$$\mathbf{C}\dot{\boldsymbol{\theta}} = -\mathbf{A}^T \mathbf{G} \mathbf{A} \boldsymbol{\theta} + \mathbf{A}^T \mathbf{G} \mathbf{b} + \mathbf{f} \quad \begin{bmatrix} \mathbf{G}^{-1} & \mathbf{A} \\ -\mathbf{A}^T & s\mathbf{C} \end{bmatrix} \begin{bmatrix} \mathbf{q} \\ \boldsymbol{\theta} \end{bmatrix} = \begin{bmatrix} \mathbf{b} \\ \mathbf{f} \end{bmatrix}$$

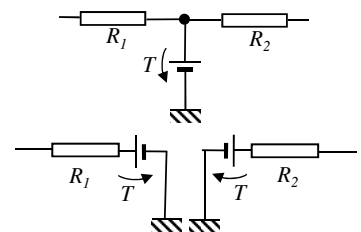
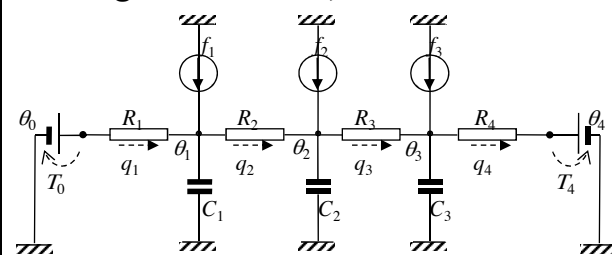
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- Plug the sources, chose a direction for fluxes



- Describe the circuit

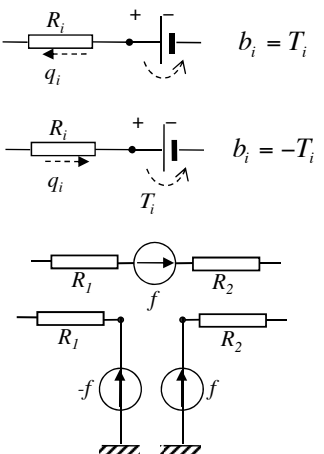
$$\boldsymbol{\theta} = [\theta_1 \quad \theta_2 \quad \theta_3]^T$$

$$\mathbf{q} = \begin{bmatrix} q_1 \\ q_2 \\ q_3 \\ q_4 \end{bmatrix} \quad \mathbf{A} = \begin{bmatrix} 1 & 0 & 0 \\ -1 & 1 & 0 \\ 0 & -1 & 1 \\ 0 & 0 & -1 \end{bmatrix} \quad \mathbf{G} = \begin{bmatrix} G_1 & 0 & 0 & 0 \\ 0 & G_2 & 0 & 0 \\ 0 & 0 & G_3 & 0 \\ 0 & 0 & 0 & G_4 \end{bmatrix} \quad \mathbf{b} = \begin{bmatrix} T_0 \\ 0 \\ 0 \\ -T_4 \end{bmatrix}$$

$$\mathbf{C} = \begin{bmatrix} C_1 & 0 & 0 \\ 0 & C_2 & 0 \\ 0 & 0 & C_3 \end{bmatrix} \quad \mathbf{f} = [f_1 \quad f_2 \quad f_3]^T$$

- Give the solution

$$\mathbf{C}\dot{\boldsymbol{\theta}} = -\mathbf{A}^T \mathbf{G} \mathbf{A} \boldsymbol{\theta} + \mathbf{f}_T ; \quad \text{where } \mathbf{f}_T = \mathbf{A}^T \mathbf{G} \cdot \mathbf{b} + \mathbf{f}$$



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- state-space model, if $\exists C^{-1}$

$$\dot{\theta} = A_s \theta + B_s f_T$$

$$y = C_s \theta + D_s f_T$$

where θ state vector; contains the temperatures

$f_T = [b \ f]^T$ - input vector; contains sources

y - output vector; variables of interest,

$A_s = -C^{-1} A^T G A$ dynamic matrix, if $\exists C^{-1}$

$B_s = C^{-1} [A^T G \ I]$ - command matrix, if $\exists C^{-1}$

C_s - observation matrix

D_s - feed-through matrix

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- Euler approximation

$$\left. \frac{\partial \theta}{\partial t} \right|_{m,n} \cong \frac{\theta_{m,n}^{p+1} - \theta_{m,n}^p}{\Delta t}$$

- evaluation at the previous time-step

$$\dot{\theta} = A_s \theta + B_s f_T$$

$$\frac{\theta_{p+1} - \theta_p}{\Delta t} = A_s \theta_p + B_s f_{T,p}$$

$$\theta_{p+1} = (I + \Delta t A_s) \theta_p + \Delta t B_s f_{T,p}$$

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• Stability

- choice of the spatial discretization
- imposes the time-step

• 1D :

$$\frac{\alpha \Delta t}{(\Delta x)^2} \leq \frac{1}{2}$$

2D :

$$\frac{\alpha \Delta t}{(\Delta x)^2} \leq \frac{1}{4}$$

$$\alpha = \lambda / (\rho c) \quad \text{thermal diffusivity}$$

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• Euler approximation

$$\left. \frac{\partial \theta}{\partial t} \right|_{m,n} \cong \frac{\theta_{m,n}^{p+1} - \theta_{m,n}^p}{\Delta t}$$

• evaluation at the same time-step

$$\dot{\theta} = \mathbf{A}_s \theta + \mathbf{B}_s \mathbf{f}_T$$

$$\frac{\theta_{p+1} - \theta_p}{\Delta t} = \mathbf{A}_s \theta_{p+1} + \mathbf{B}_s \mathbf{f}_{T,p}$$

$$\theta_{p+1} = (\mathbf{I} - \Delta t \mathbf{A})^{-1} (\theta_p + \Delta t \mathbf{B}_s \mathbf{f}_{T,p})$$

Convection

Constitutive law: Newton

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Heat transfer in fluids

- *advection*
- *diffusion*

Steady-state

Fluid – solid wall

- heat flux at the wall

$$\varphi_p = -\lambda_{ps} \frac{\partial \theta}{\partial n} \Big|_{ps} = -\lambda_{pf} \frac{\partial \theta}{\partial n} \Big|_{pf}$$

- temperature continuity at the wall

$$\theta_s \Big|_M = \theta_f \Big|_M$$

Convection

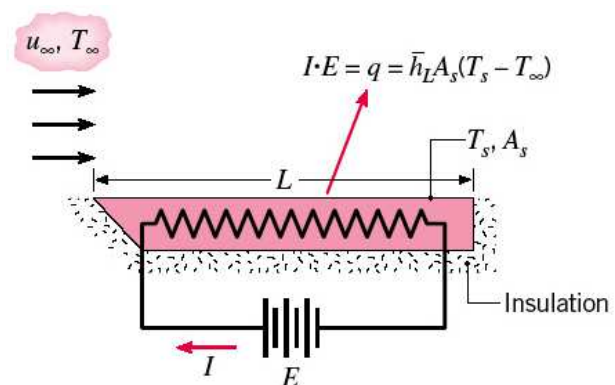
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Newton law

$$\varphi_p = h_c (\theta_p - \theta_M)$$



Convection problem

- equivalent temperature
- convective heat coefficient

Convection

Constitutive law: Newton

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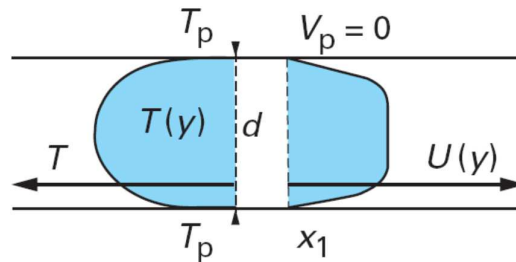
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• Internal flow



$$\dot{m}_x = \int_0^d \rho v dy = d \rho v_M$$

$$v_M \equiv \frac{1}{d} \int_0^d v dy$$

$$\dot{H}_x = \int_0^d \rho c v \theta dy$$

$$\theta_M = \frac{\dot{H}_x}{c \dot{m}_x} = \frac{\int_0^d \rho v \theta dy}{\int_0^d \rho v dy}$$

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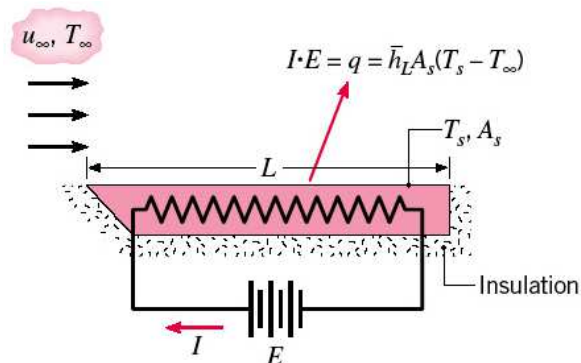
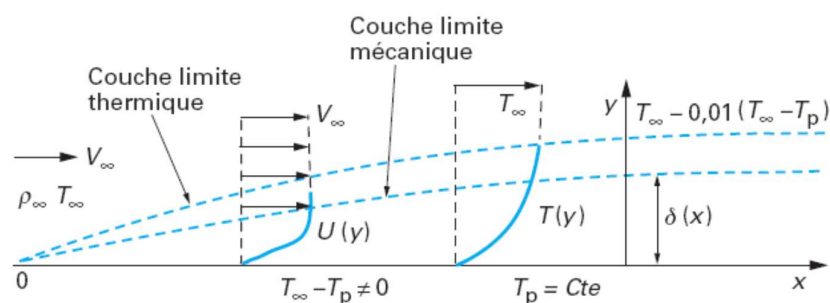
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• External flow

$$\theta_M = \theta_\infty$$



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• Approximative values

	Convection	
	natural (W/m ² K)	forced (W/m ² K)
gaz	3 ...30	12 ... 200
liquid	30 ...300	200...7500

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• Mathematical expression of a physical law

$$f(E_1, \dots, E_p) = 0$$

Theorem of Vachy – Buckingham

$$F(\pi_1, \dots, \pi_{p-q}) = 0$$

q number of fundamental units
($L, T, M, I, \dots$)

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• Example : heat exchange in an infinite pipe

Values which influence the thermal flux in the pipe

	Value	Symbol	SI unit	Dimension
1	Pipe diameter	(a) D	m	$[L]$
2	Fluid velocity	(b) U_{∞}	m/s	$[LT^{-1}]$
3	Fluid density	(c) ρ	kg/m ³	$[ML^{-3}]$
4	Fluid viscosity	(d) μ	kg/(m·s)	$[ML^{-1}T^{-1}]$
5	Thermal conductivity of the fluid	(e) λ	W/(m·K)	$[MLT^{-3}\theta^{-1}]$
6	Specific heat at constant pressure	(f) c_p	J/(kg·K)	$[L^2T^{-2}\theta^{-1}]$
7	Heat convection coefficient	(g) h_c	W/(m ² K)	$[MT^{-3}\theta^{-1}]$

7 [physical variables] - 4 [dimensions (L, M, T, θ)] = 3 [products]

$$\pi = D^a \cdot v^b \cdot \rho^c \cdot \mu^d \cdot \lambda^e \cdot c_p^f \cdot h_c^g$$

dimensions :

$$[\pi] = [L]^{a+b-3c-d+e+2f} [M]^{c+d+e+g} [T]^{-b-d-3e-2f-3g} [\theta]^{-e-f-g}$$

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– adimensional products (4 eqs. 7 unknowns)

$$\begin{cases} a + b - 3c - d + e + 2f = 0 \\ c + d + e + g = 0 \\ -b - d - 3e - 2f - 3g = 0 \\ -e - f - g = 0 \end{cases}$$

– a possible solution

$$a = 1; e = -1; g = 1; b = c = d = f = 0; \Rightarrow \pi_1 = \frac{h \cdot D}{\lambda} \equiv Nu$$

$$a = 1; b = 1; c = 1; d = -1; e = f = g = 0; \Rightarrow \pi_2 = \frac{\rho \cdot v \cdot D}{\mu} \equiv Re$$

$$d = 1; e = -1; f = 1; a = b = c = g = 0; \Rightarrow \pi_3 = \frac{\mu \cdot c_p}{\lambda} \equiv Pr$$

$$F(Nu, Re, Pr) = 0$$

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- **Reynolds** : ratio between inertial and viscous forces

$$Re = \frac{\rho v D}{\mu} = \frac{\rho v^2 D^2}{\mu v D}$$

- *inertia*

$$\rho v^2 D^2 = \rho v D^2 \cdot v = \rho l D^2 \cdot \frac{v}{t} \propto \rho V \cdot \frac{v}{t} = m \cdot a$$

- *forces viscous*

$$\mu v D \propto \mu S \frac{dv}{dz}$$



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- **Prandtl** : viscous diffusion rate / thermal diffusion rate

$$Pr = \frac{\mu \cdot c_p}{\lambda} = \frac{\mu / \rho}{\lambda / (\rho \cdot c_p)} = \frac{\nu}{\alpha}$$

- **Nusselt** : convective heat transfer / conductive heat transfer

$$Nu = \frac{h D}{\lambda} = \frac{h \Delta T}{\lambda \frac{\Delta T}{D}} = \frac{D / \lambda}{1 / h}$$

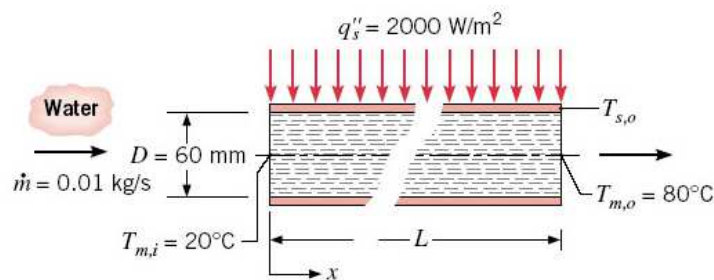
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• inside flow



h_c depends on:

- v_M , mean velocity of the fluid, m/s
- ρ , density of fluid, kg/m³
- c , specific heat of fluid, J/kg °C
- μ , dynamic viscosity of fluid, Pa s
- λ , thermal conductivity of fluid, W/m °C
- D , inner diameter of pipe, m
- x , abscissa, m.

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• internal flow

Four fundamental units (four adimensional numbers):

$$Nu = \frac{h D}{\lambda} \text{ Nusselt, heat exchange between fluid and wall;}$$

$$Re = \frac{\rho v_m D}{\mu} \text{ Reynolds, flow type:}$$

$Re < 2000$ laminar flow
 $Re > 3000$ turbulent flow

$$Pr = \frac{\mu c}{\lambda} \text{ Prandtl, thermal properties of the fluid}$$

$$\frac{x}{D} \text{ adimensional abscissa}$$

$$Nu = 0.023 Re^{0.8} Pr^{0.33}$$

Convection

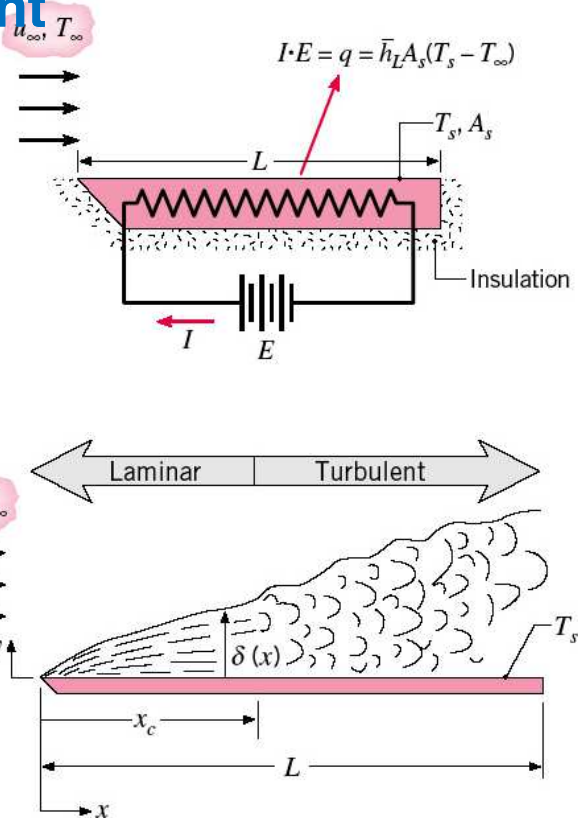
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external flow

$$\overline{Nu} = 0.035 Re^{0.80} Pr^{0.33}$$



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- Gravity and buoyancy
- Grashof number: equivalent to Reynolds number for natural convection

$$Gr \equiv \frac{g \rho \beta (\theta_p - \theta_\infty) l^3}{\mu^2}$$

- buoyancy (Archimedes)
- viscous forces

Correlation for natural convection

$$Nu = C(Gr Pr)^n$$

$$Ra \equiv Gr Pr \text{ (Rayleigh)}$$