

M8: Building design: multidisciplinary approach

Hydraulics

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Hydraulics-pressurized flows

Scope of the course/project

- Design and size a water distribution network (water supply/heating networks)
- Calculate flow rate distribution in a water supply network with respect of water demand and operating pressures

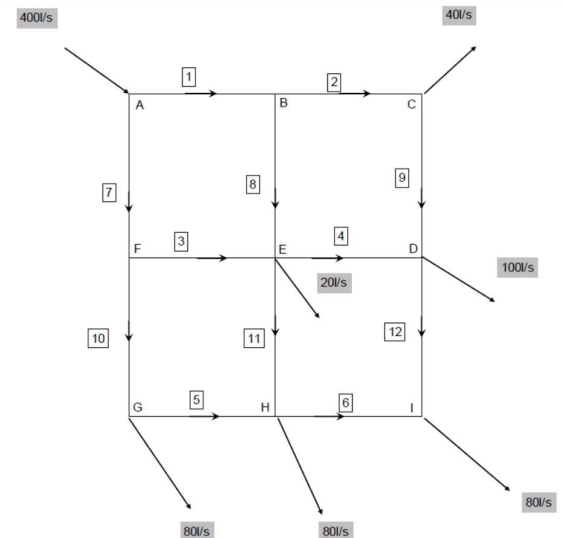
... develop the ability to implement design methods into algorithms (using Matlab)

Course overview

-Scope

-Reminders and exercise (~2-3h)

➔ Calculate the flow distributions in a meshed-type network



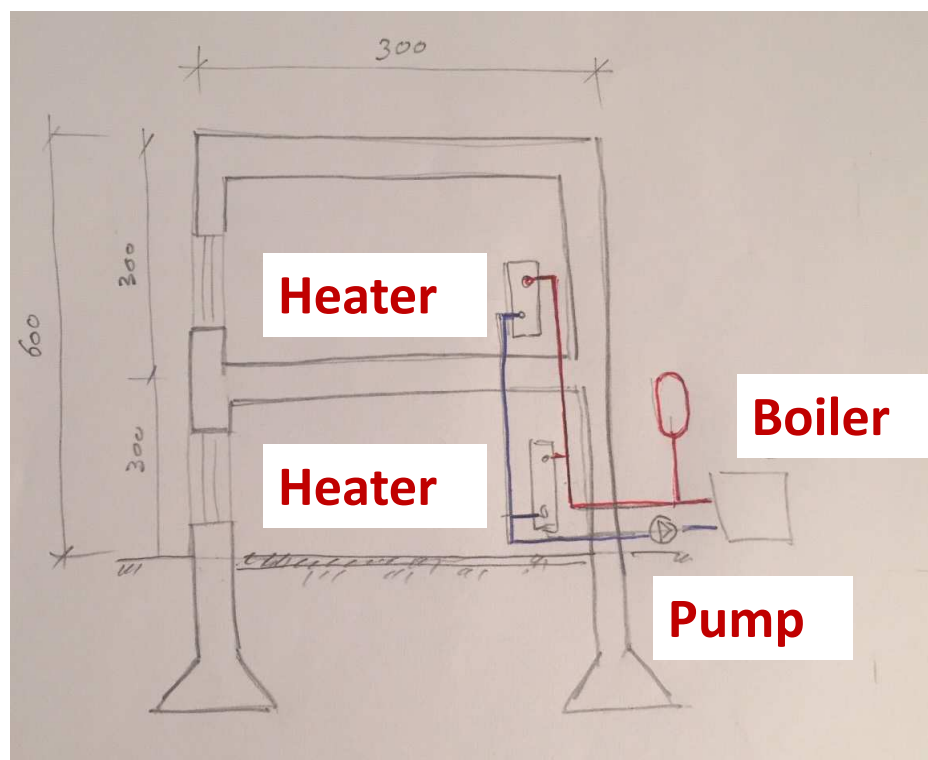
-Project

Design and size a water distribution network for the heating (2 storey building)

3

Project Scope

Design and size a water distribution network for the heating of two storey-building (pump, pipes...) according to the heating demand



4

Reminders

- Fluids properties
- Conservation equations
- Head losses
- Pipe networks
- Centrifuge pump and operating point of a system

5

Reminders: fluids properties

Fluid/solid

React differently to shear stress: solid deform only a finite amount, whereas fluids deform continuously until the stress is removed

Liquid/gas

Both liquids and gases are fluids: differ strongly in both degree of compressibility and formation of a free surface (interface) for the liquid

Liquids: incompressible, gases range from compressible to incompressible

Density (at standard conditions: 20°C and 101.325 kPa-sea level atmospheric pressure):

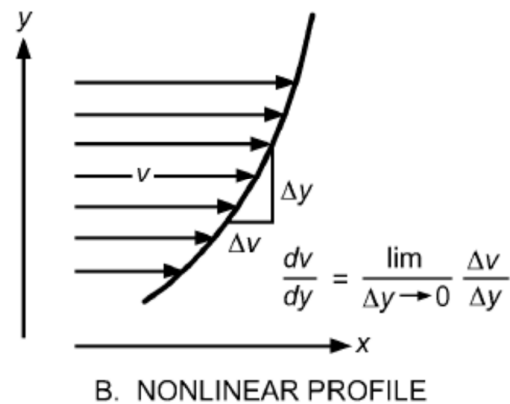
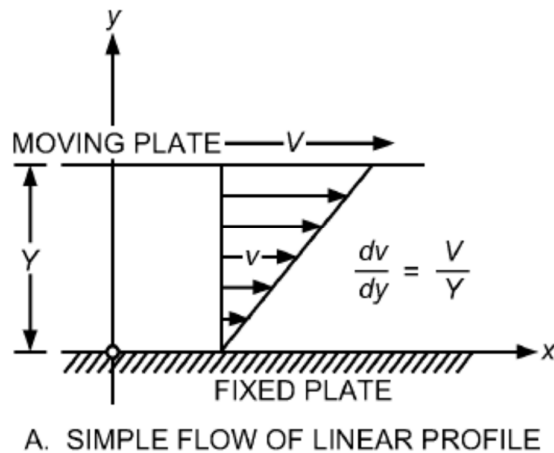
$$\rho_{\text{water}} = 998 \text{ kg/m}^3 \quad \rho_{\text{air}} = 1.21 \text{ kg/m}^3$$

6

Reminders: fluids properties

Viscosity: Resistance of adjacent fluid layers to shear.

μ is the absolute or dynamic viscosity of the fluid.



In **Newtonian fluids**, the rate of deformation proportional to the shearing stress

$$F/A = \mu (V/Y)$$

In **non Newtonian fluids**, velocity and shear stress may vary across the flow field

$$\tau = \mu \frac{dv}{dy}$$

Reminders: fluids properties

Viscosity: Resistance of adjacent fluid layers to shear.

μ is the absolute or dynamic viscosity of the fluid.

Absolute viscosity has dimensions of force $\times$ time/length²

$$\mu_{\text{water}} = 1.01 \cdot 10^{-3} \text{ N.s/m}^2$$

$$\mu_{\text{air}} = 18.1 \cdot 10^{-6} \text{ N.s/m}^2$$

(At standard indoor conditions)

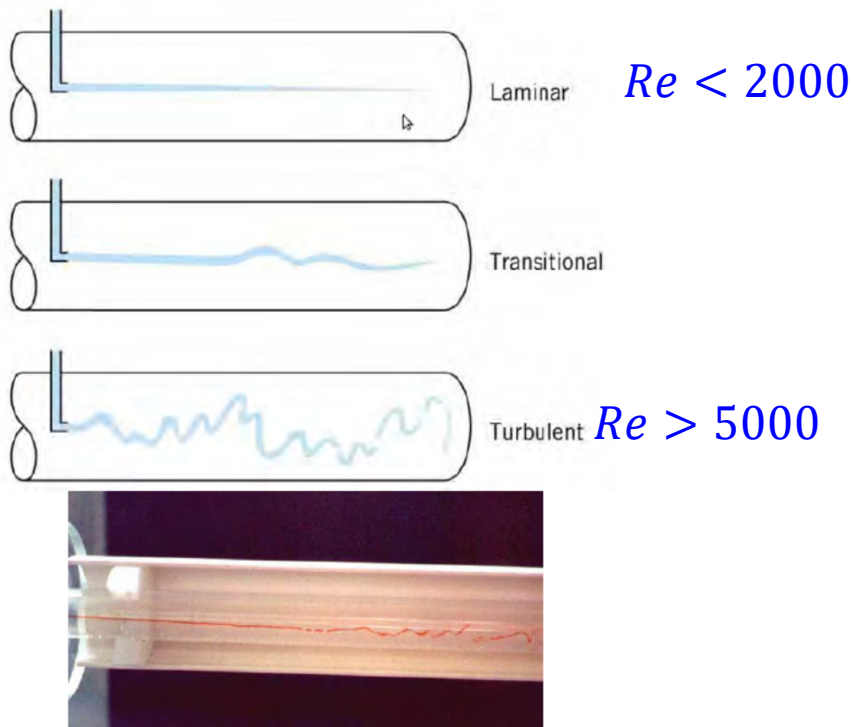
***Kinematic viscosity** ν is sometimes used in lieu of dynamic viscosity μ

$$\nu = \mu / \rho$$

$$\nu_{\text{water}} = 1.01 \cdot 10^{-6} \text{ m}^2/\text{s}^{-1}$$

$$\nu_{\text{air}} = 15. \cdot 10^{-6} \text{ m}^2/\text{s}^{-1}$$

Reminders: fluids properties



Reynolds number:

$$Re = \frac{\text{Inertial forces}}{\text{Viscous forces}}$$

(dimensionless ratio)

$$Re = \frac{VL}{\nu}$$

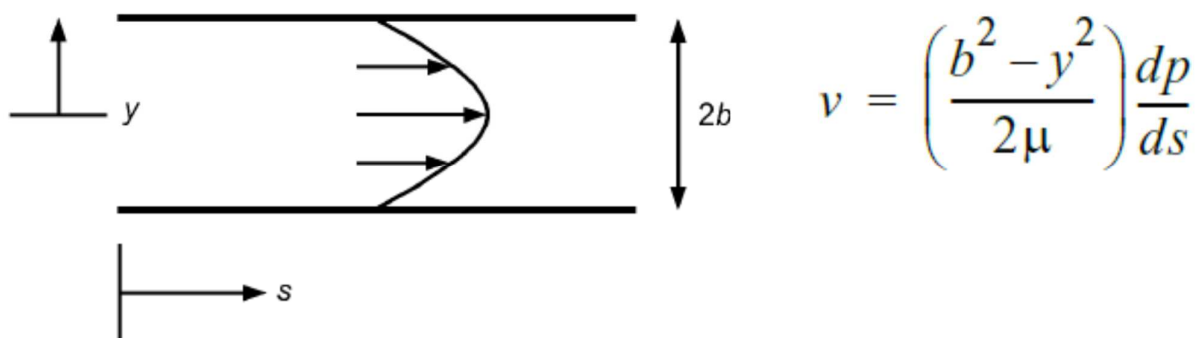
L is a characteristic length scale

ν is the kinematic viscosity of the fluid.

Reminders: fluids properties

Laminar vs turbulence flows

For steady, fully developed **laminar flow** between 2 parallel plates:



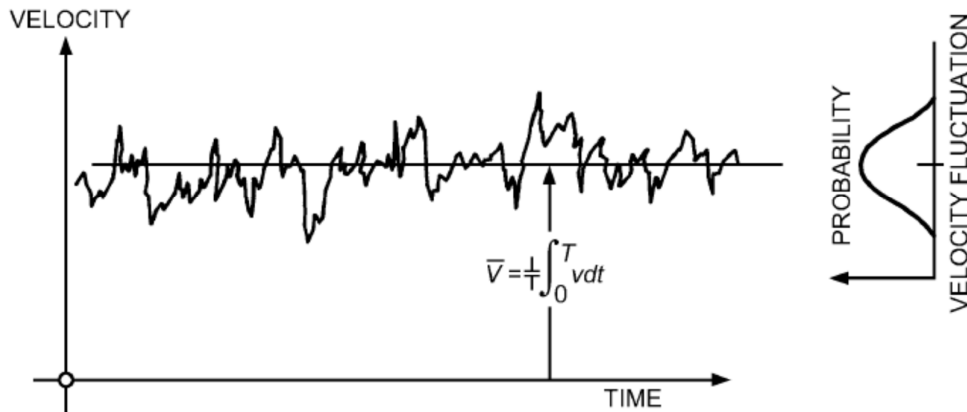
The resulting parabolic velocity profile in a wide rectangular channel is commonly called **Poiseuille flow**.

Maximum velocity occurs at the centerline ($y = 0$), and the average velocity U is $2/3$ of the maximum velocity.

Reminders: fluids properties

Laminar vs turbulence flows

Fluid flows are generally **turbulent**, involving random perturbations or fluctuations of the flow (velocity and pressure):



Turbulence can be quantified statistically.

The velocity used is the **time-averaged velocity**.

In flow through pipes, tubes, and ducts, the **characteristic length scale** is the hydraulic diameter D_h , given by

$$D_h = 4A/P_w$$

A: cross-sectional area of the pipe, P_w : wetted perimeter

11

Reminders: mass conservation

Continuity equation

$$\dot{m} = \int \rho v dA = \text{constant}$$

where $\dot{m}$ is the mass flow rate (kg/s) across the area normal to flow, V is fluid velocity normal to differential area dA , and ρ is fluid density.

When flow is incompressible ($\rho = \text{constant}$) in a pipe:

$$\dot{m} = \rho VA \quad Q = \dot{m}/\rho = U \times A = \text{cst}$$

Where U , the average velocity is defined as follows:

$$U = \frac{Q}{A} = \frac{1}{A} \iint_S \vec{V} \cdot \vec{n} dA$$

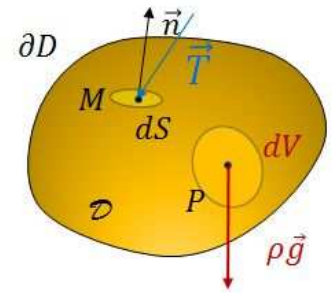
Q is **volumetric flow rate** (m^3/s)

12

Reminders: momentum conservation

D: material domain

$\vec{n}$ the external normal



Momentum principle :

$$\sum \vec{F}_{ext} = \frac{d(m\vec{V})}{dt}$$

$m\vec{V}$ is the momentum.

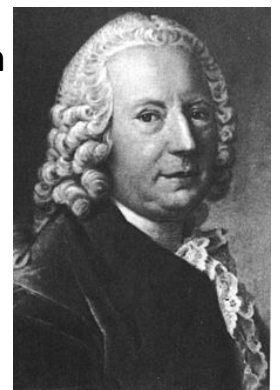
with $m = \iiint_D \rho dV$

13

Reminders: energy conservation (Bernoulli equation)

The Bernoulli equation derives from the **conservation of momentum** and express the **conservation of energy** along a streamline;

The change in energy content ΔE (per unit mass) of flowing fluid is a result of the work (per unit mass) done on the system (heat absorbed or rejected neglected).



Fluid energy is composed of kinetic, potential (elevation z , pressure) energies. Per unit mass of fluid, the relation of the energy change between two sections of the system is



1700-1782

$$\Delta E_{mechanical} = \Delta \left(gz + \frac{P}{\rho} + \frac{1}{2} V^2 \right) = E_{ext}$$

E_{ext} is the external work from a fluid machine (positive for a pump or blower)

V is the velocity along a streamline

14

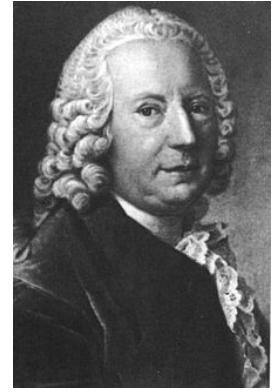
Reminders: energy conservation (Bernoulli equation)

If there is no viscous frictional forces, no external work ($E_{\text{ext}}=0$), no heat transfer, it comes:

$$gz + \frac{P}{\rho} + \frac{1}{2}V^2 = B = \text{cst}$$

B is the Bernoulli constant

Energy (per unit mass) is conserved along a streamline.



1700-1782

Alternative form is:

$$z + \frac{P}{\rho g} + \frac{1}{2g}V^2 = \frac{B}{g} = H$$

Energy per unit weight (**Head**)

15

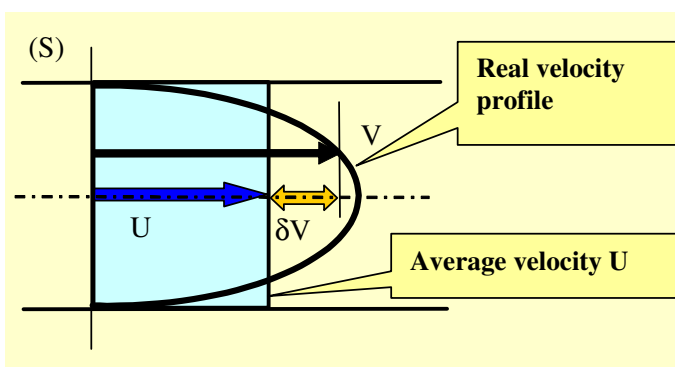
Reminders: energy conservation (Bernoulli equation)

When the **velocity V varies significantly across the cross section**, the kinetic energy term has to be correctly calculated:

$$z + \frac{P}{\rho g} + \frac{\alpha}{2g}U^2 = H$$

The velocity used is the **average velocity (U)**, so that a **kinetic energy factor α** ($\alpha > 1$) has to be calculated to take into account true kinetic energy of the velocity profile.

α expresses the ratio of the true kinetic energy of the velocity profile to kinetic energy of the average velocity :



For laminar flow in a pipe, $\alpha = 2$
For turbulent flow, $\alpha \approx 1$
(cylindrical pipes)

16

Reminders: energy conservation (Bernoulli equation)

Real fluid flow

If there are now **viscous frictional forces** and **external work (pump)**, conversion of mechanical energy to internal energy may be expressed as a loss ΔE_L .

The change in the Bernoulli equation between stations 1 and 2 along the conduit can be expressed as

$$\left(gz + \frac{P}{\rho} + \frac{\alpha}{2} U^2 \right)_1 + E_{pump} - \Delta E_L = \left(gz + \frac{P}{\rho} + \frac{\alpha}{2} U^2 \right)_2$$

Alternative form is (energy by unit weight):

$$\left(z + \frac{P}{\rho g} + \frac{\alpha}{2g} U^2 \right)_1 + H_{pump} - \Delta H_L = \left(z + \frac{P}{\rho g} + \frac{\alpha}{2g} U^2 \right)_2$$

ΔH_L is the head loss through friction

17

Reminders: Frictional head losses

In real-fluid flow, **a frictional shear occurs at bounding walls**, gradually influencing flow further away from the boundary. A lateral velocity profile is produced and **flow energy is converted into heat** (fluid internal energy), which is generally **unrecoverable (a loss)**. This loss in fully developed conduit flow is evaluated using the **Darcy-Weisbach equation**:

$$\Delta H_L = f \frac{L}{D} \frac{U^2}{2g}$$

where L is the length of conduit of diameter D and **f** is the **Darcy-Weisbach friction factor**.

18

Frictional head losses (Determination of friction coefficient f)

- **Phase 1: Poiseuille** equation for a laminar regime with $Re < 2000$, with a slope

$$f = \frac{64}{Re}$$

- **Phase 3: Blasius law ($Re < 10^5$):** $f = \frac{0.316}{Re^{1/4}}$

ϵ/D = Relative roughness

Hydraulically smooth regime

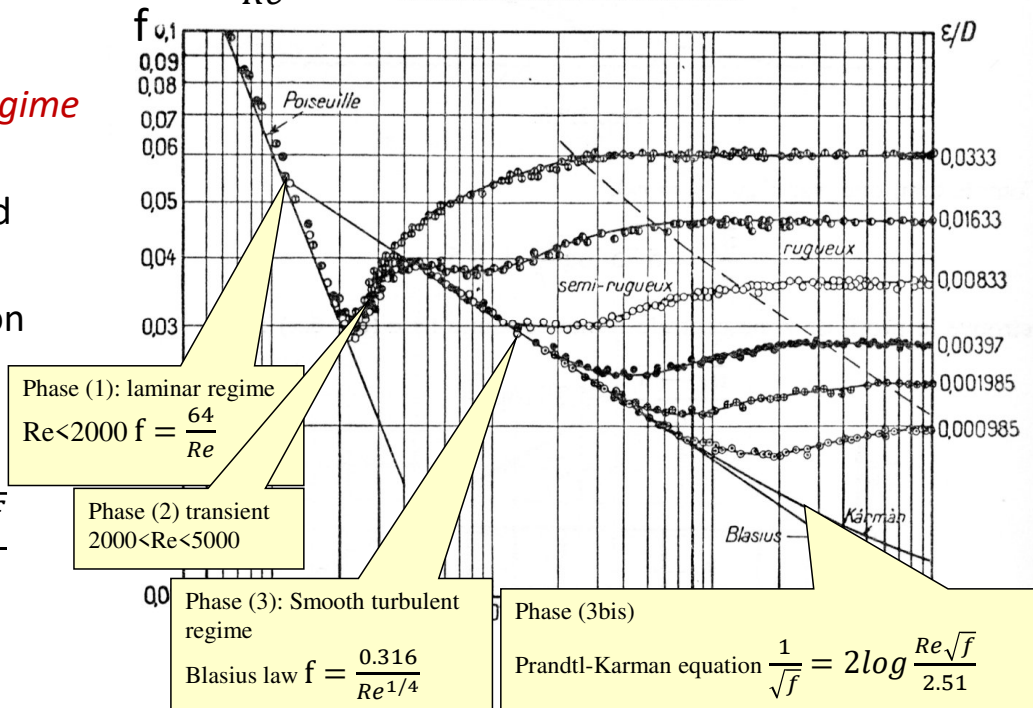
For the smallest ϵ/D and

$Re > 10^5$:

Prandtl-Karman equation better fits (phase 3bis)

Phase (3bis)

$$\frac{1}{\sqrt{f}} = 2 \log \frac{Re \sqrt{f}}{2.51}$$



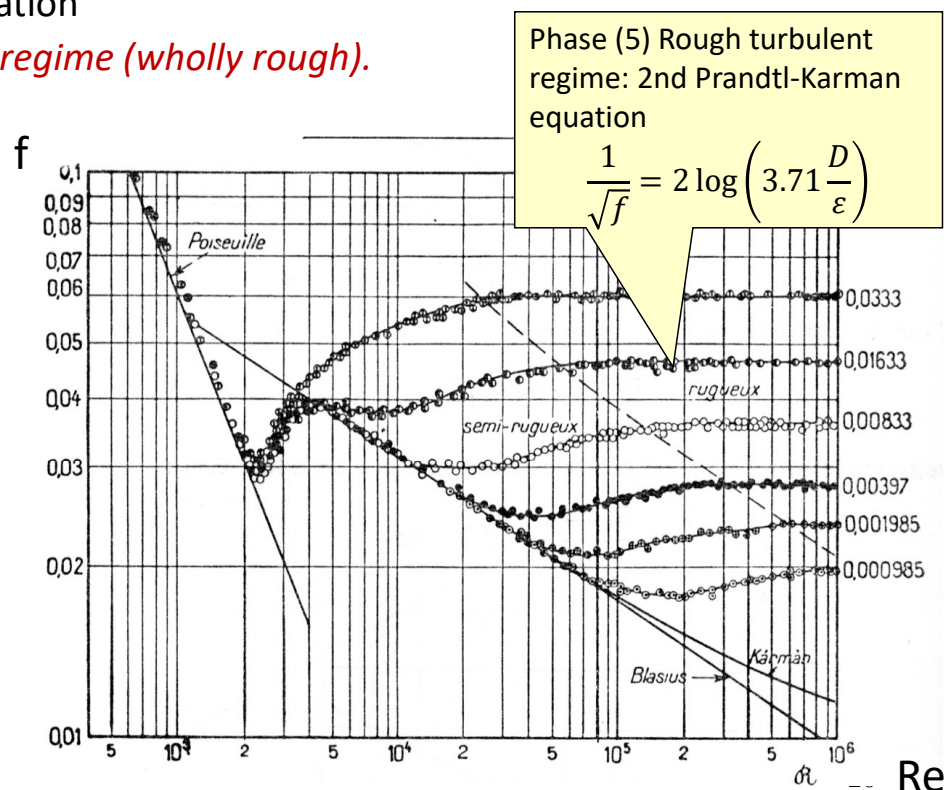
Frictional head losses (Determination of friction coefficient f)

- **Phase 5**, friction factor depends only on **relative roughness ϵ/D** ;

2nd Prandtl-Karman equation

Hydraulically rough flow regime (wholly rough).

$$\frac{1}{\sqrt{f}} = 2 \log 3.71 \frac{D}{\epsilon}$$

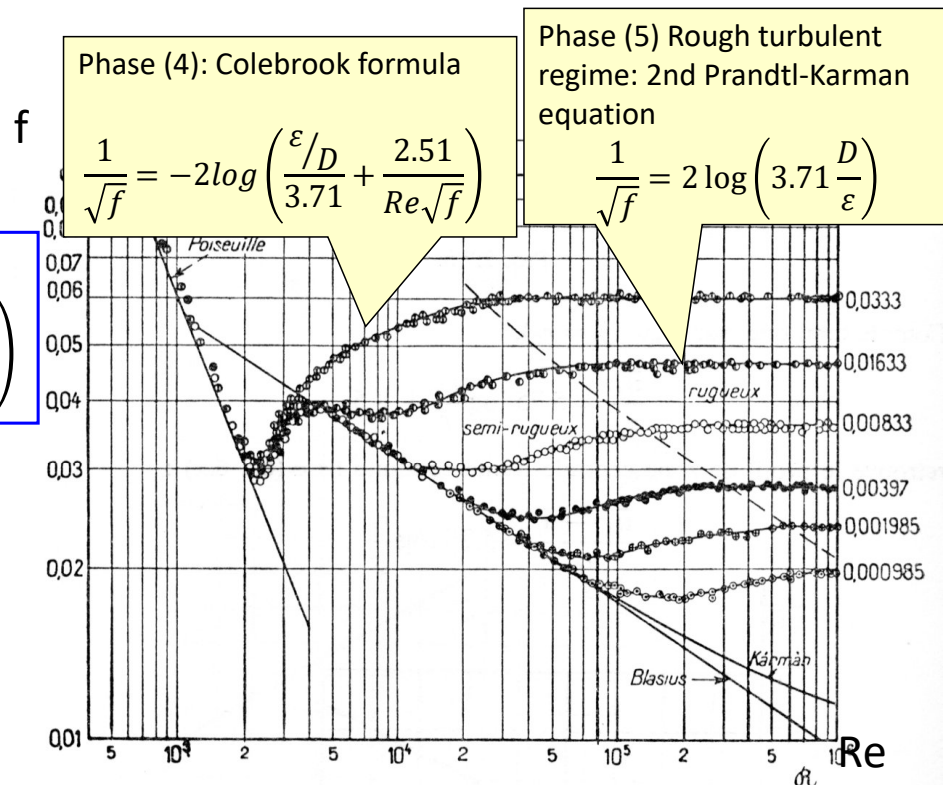


Frictionnal head losses (Determination of friction coefficient f)

Phase 4: friction coefficient depends both on Re and ϵ/D .

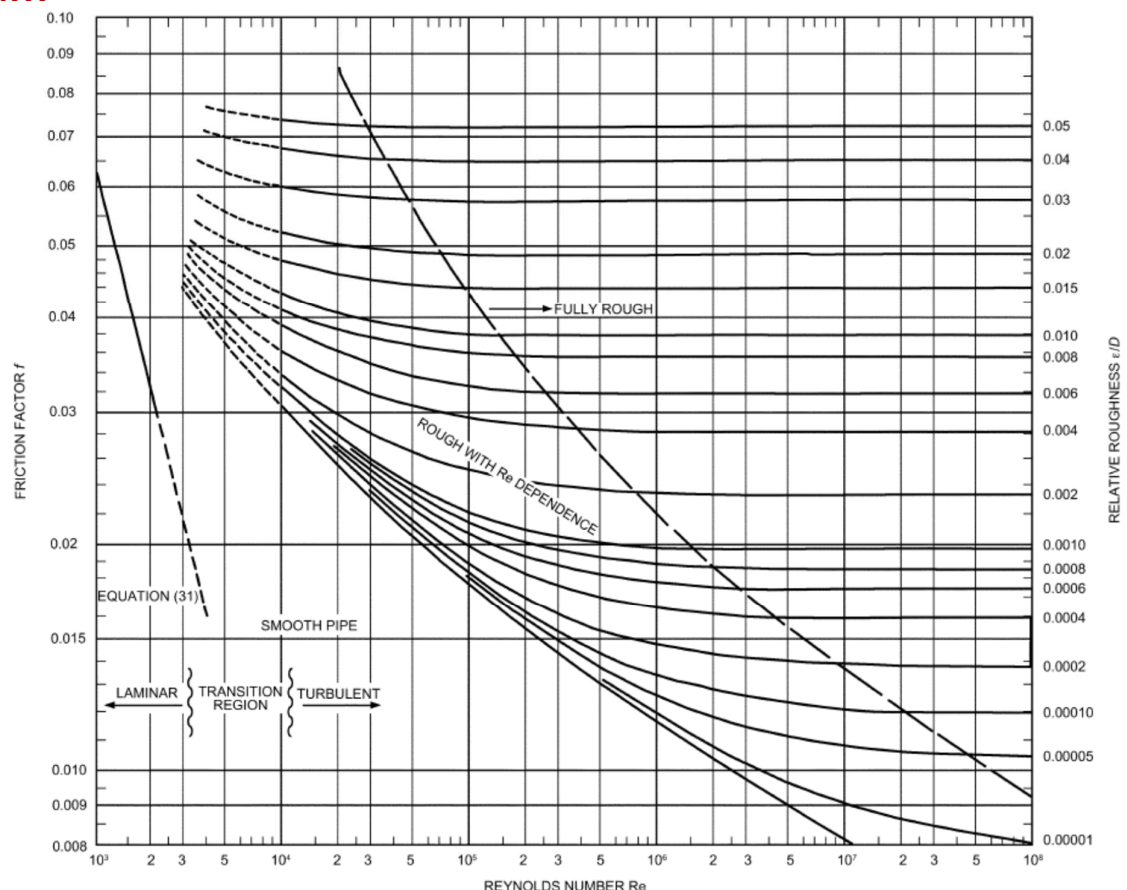
Colebrook formula

$$\frac{1}{\sqrt{f}} = -2 \log \left(\frac{\epsilon/D}{3.71} + \frac{2.51}{Re \sqrt{f}} \right)$$



Frictionnal head losses (Determination of friction coefficient f)

Moody diagram



Frictionnal head losses (Industrial pipes)

Other empirical formulae (not exhaustive)

Authors and conduit type	empirical formulae notation: $U(\text{m/s}), D(\text{m}), Q(\text{m}^3/\text{s})$ $j = \frac{\Delta H}{L}$	Comments
Chézy cast iron	$U = C \sqrt{\frac{D \cdot j}{4}}$ then $\Delta H = \frac{U^2 L}{2 \cdot g \cdot D} \left(\frac{8 \cdot g}{C^2} \right)$	$40 < C < 100$ then $0,008 < \frac{8 \cdot g}{C^2} < 0,022$
Darcy cast iron $0,05 < D(\text{m}) < 0,5$	$\frac{D \cdot j}{4} = \left(\alpha + \frac{\beta}{D} \right) \cdot U^2$ then $\Delta H = \frac{U^2 L}{2 \cdot g \cdot D} \left(8 \cdot g \left(\alpha + \frac{\beta}{D} \right) \right)$	$\alpha = 0,000507$ $\beta = 0,00001294$ then $0,04 < 8 \cdot g \left(\alpha + \frac{\beta}{D} \right) < 0,06$
Lévy cast iron $0,08 < D(\text{m}) < 3$	$U = \beta \cdot \sqrt{\frac{D \cdot j}{2} \cdot \left(1 + 3 \sqrt{\frac{D}{2}} \right)}$ then $\Delta H = \frac{U^2 L}{2 \cdot g \cdot D} \left(\frac{4 \cdot g}{\beta^2 \left(1 + 0,75 \sqrt{D} \right)} \right)$	$0,013 < \frac{4 \cdot g}{\beta^2 \left(1 + 0,75 \sqrt{D} \right)} < 0,078$ $\beta = 36,4$ (fonte neuve) $\beta = 20,5$ (fonte en service)
Flamant and Maning cast iron $D(\text{m}) < 1,3$	$D \cdot j = 0,00092 \cdot 4 \sqrt{\frac{U^7}{D}}$	cast iron in service

23

Reminders: Head losses/ Minor losses

Bends, entrances, exits, contractions, enlargements...

We treat the entire effect of a disturbance as if it occurs at a single point in the flow direction.

By treating these losses as a local phenomenon, they can be related to the velocity by the **loss coefficient K**:

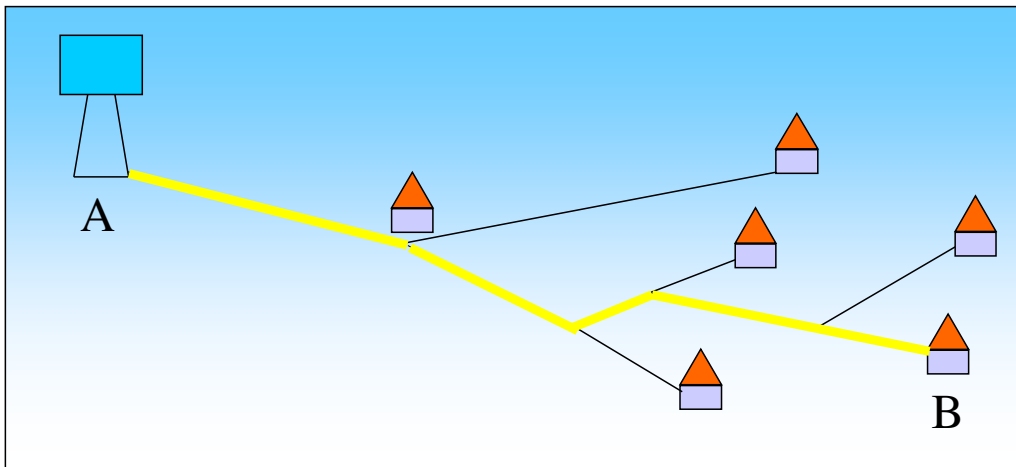
$$\Delta H = K \frac{U^2}{2g}$$

K depends on the particularity of the fittings

24

Reminders: Pipe networks

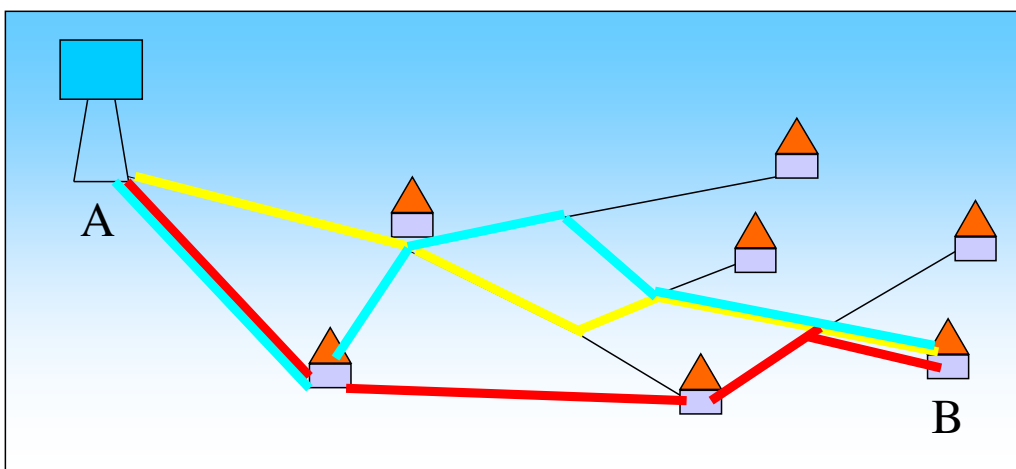
Branched network



Hydraulic characteristic	ZONE
An <u>only path</u> from A to B	Rural area

Reminders: Pipe networks

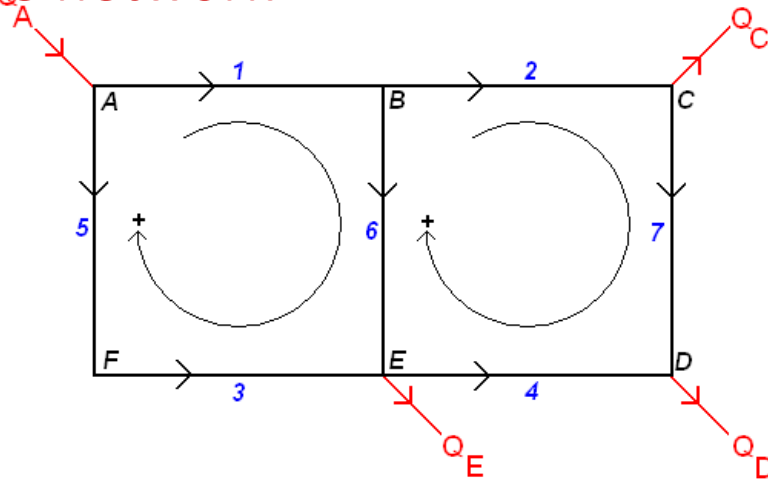
Meshed-type network (several loops)



Hydraulic characteristic	ZONE
<u>Several possible paths</u> from A to B	Urban

Reminders: Pipe networks

Meshed-type network

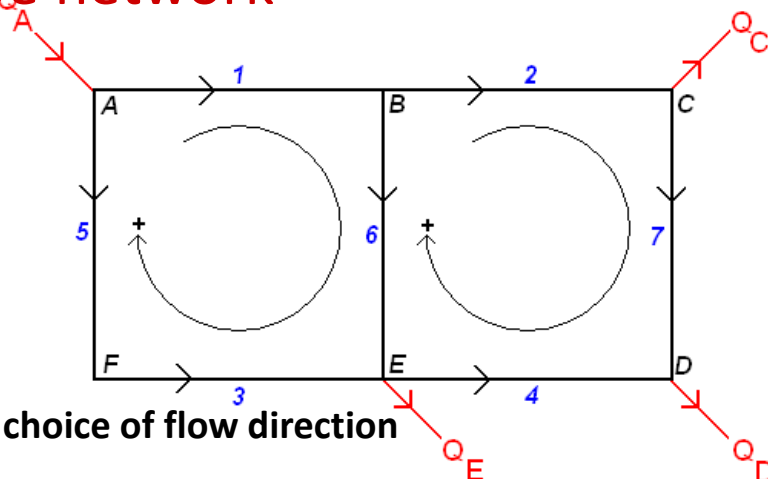


Unknown= Flow rates in sections/Nodes heads

Solution = ?

Reminders: Pipe networks

Meshed-type network



-Hypothesis on the choice of flow direction

-Arbitrary initial flow rate (Q) distribution vs arbitrary head (H) distribution

-Mass conservation for each node

-Head losses around each loop ($\sum \Delta H = 0$)

Unknown= Flow rates in sections/Nodes heads

Solution = ?

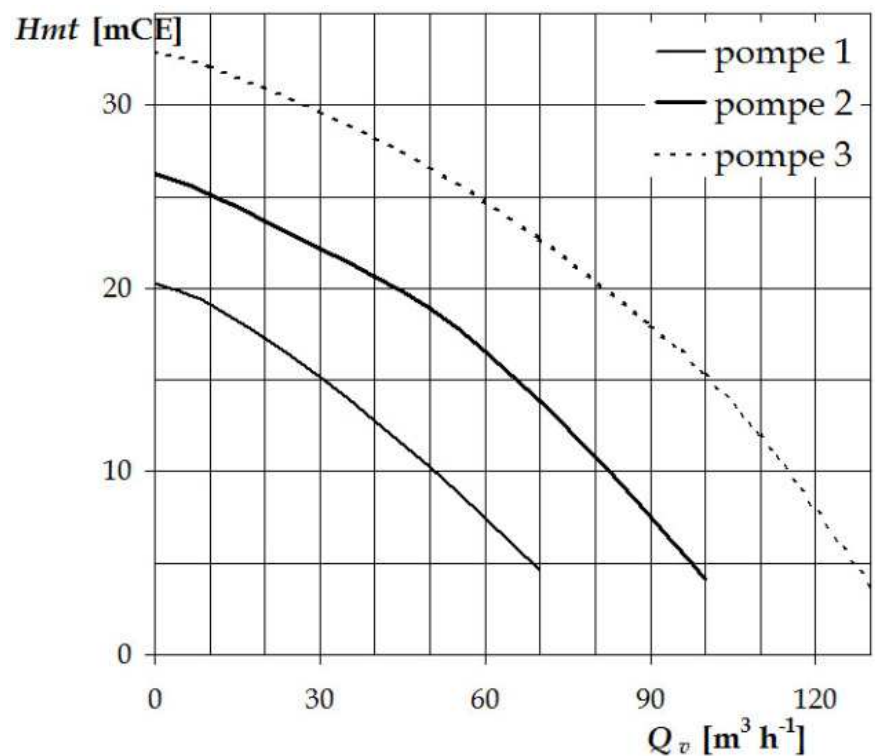
Exercise: computation of flow rates/heads in a meshed-type network

- Hypothesis on the choice of flow direction
- Arbitrary initial flow rate (Q) distribution vs Arbitrary head (H) distribution
- Mass conservation for each node
- Head losses around each loop ($\sum \Delta H = 0$)

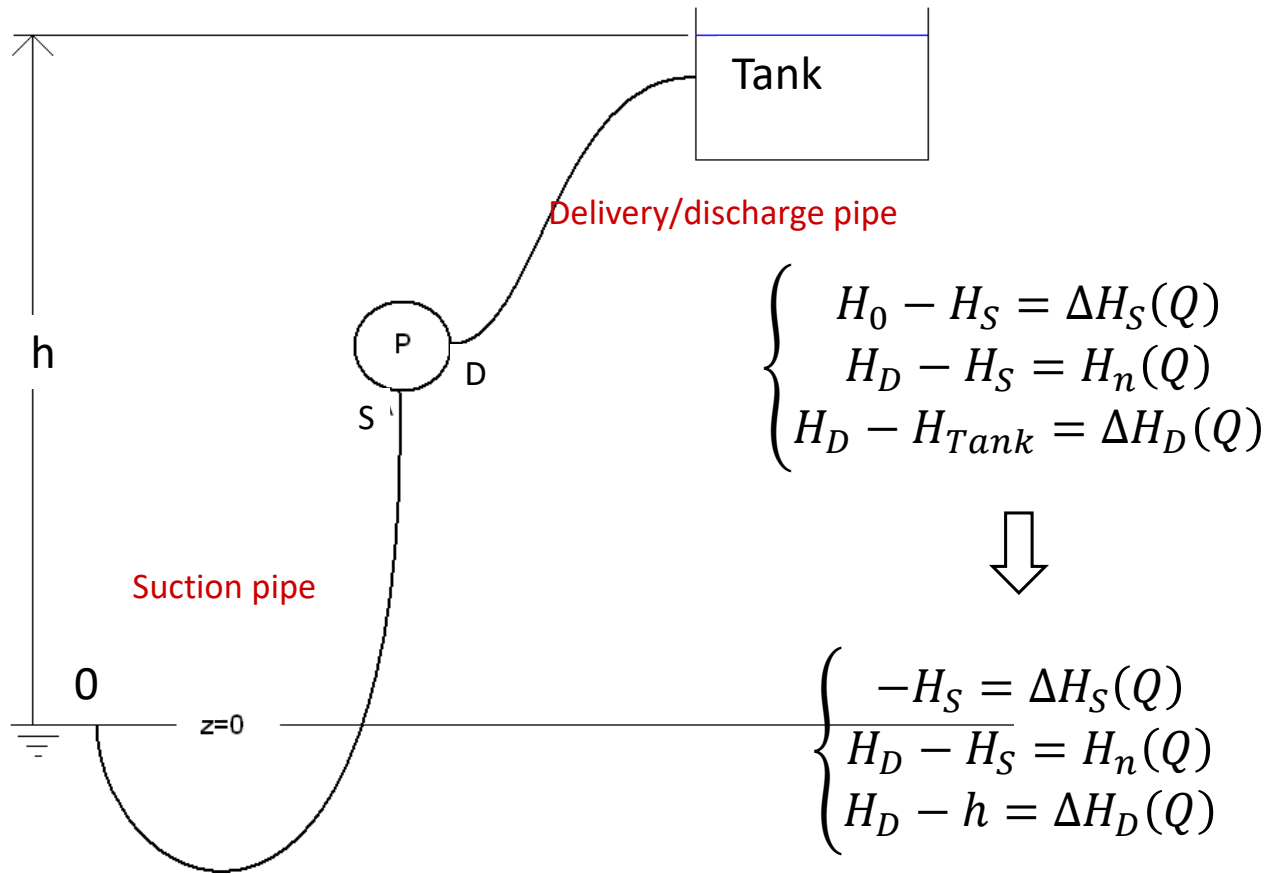
Unknown= Flow rates in section/Nodes heads

Solution = ?

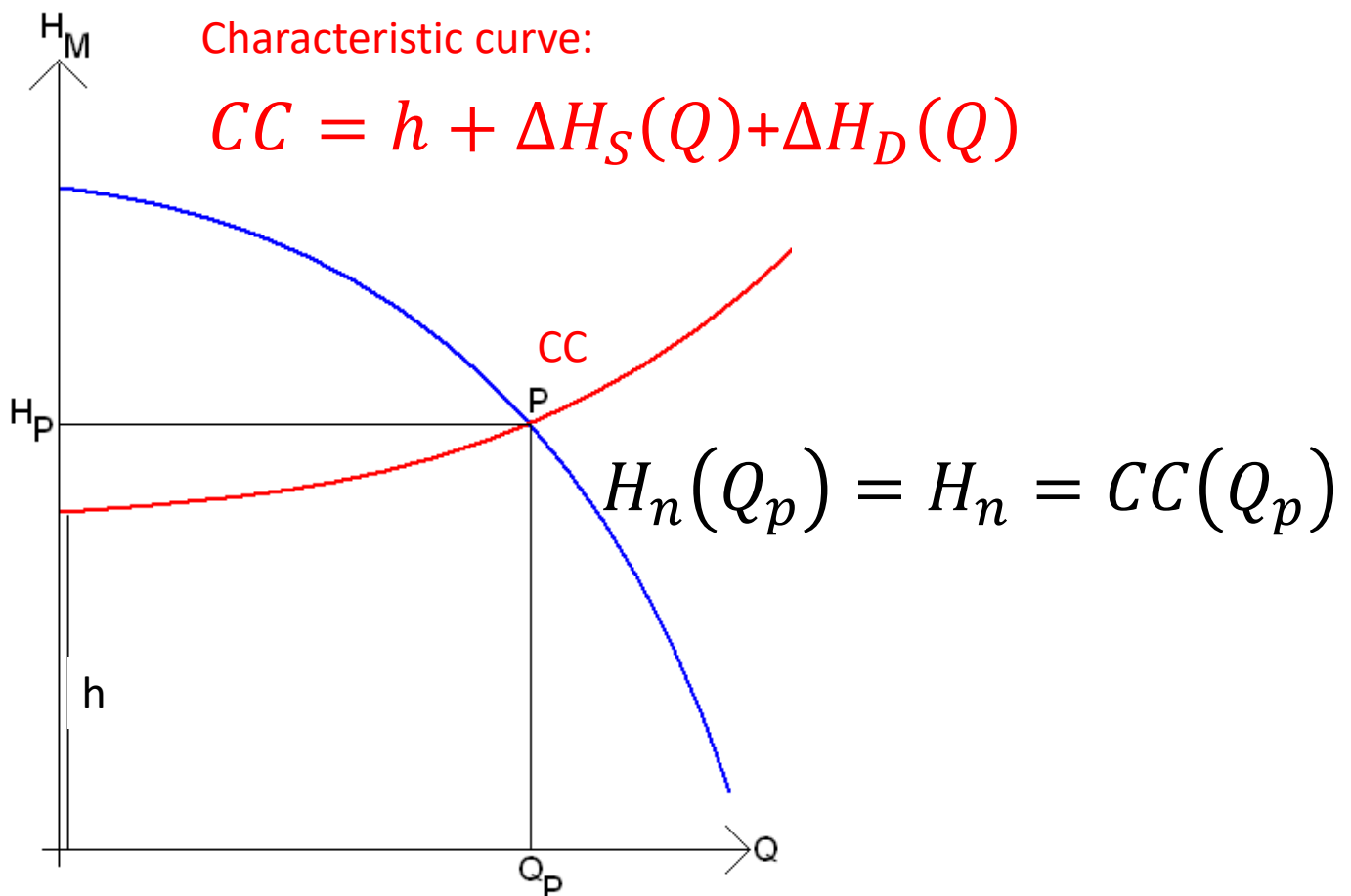
Reminders: centrifuge pump



Reminders: operating point of a system

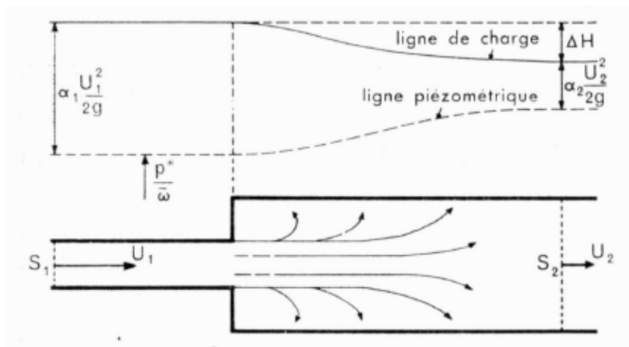


Reminders: operating point of a system



Appendix : Head losses/ Minor losses

Exemple des changements de section

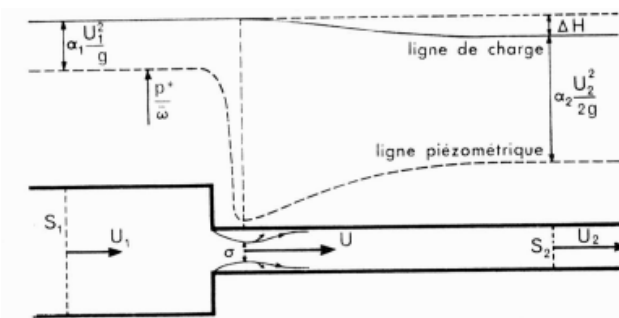


Elargissement brusque

D'après Borda

$$\Delta H = \frac{(U_1 - U_2)^2}{2g}$$

$$\Rightarrow \Delta H = \frac{U_1^2}{2g} \cdot \left(1 - \frac{U_2}{U_1}\right)^2 = \frac{U_1^2}{2g} \cdot \left(1 - \frac{S_1}{S_2}\right)^2 \text{ soit } K = \left(1 - \frac{S_1}{S_2}\right)^2$$



Rétrécissement brusque

D'après Borda

$$\Delta H = \frac{(U_1 - U_2)^2}{2g}$$

$$\Rightarrow \Delta H = \frac{U_2^2}{2g} \cdot \left(\frac{U_1}{U_2} - 1\right)^2 = \frac{U_2^2}{2g} \cdot \left(\frac{S_2}{\sigma} - 1\right)^2 \text{ soit } K = \left(\frac{1}{C_c} - 1\right)^2$$

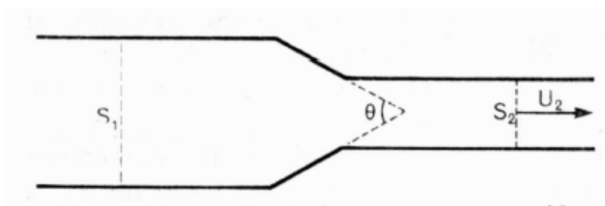
Dans un rétrécissement brusque **la perte de charge se localise en aval de la section contractée** telle que $\sigma = C_c \cdot S_2$

avec

$$C_c = 0,63 + 0,37 \cdot \left(\frac{S_2}{S_1}\right)^3 \quad 33$$

Appendix : Head losses/ Minor losses

Exemple des changements de section



Rétrécissement en biseau

$$K = \left(\frac{1}{C_c} - 1\right)^2 \cdot \sin \theta \quad \text{pour} \quad \theta < \frac{\pi}{2}$$

$$K = \left(\frac{1}{C_c} - 1\right)^2 \quad \text{pour} \quad \theta > \frac{\pi}{2}$$

avec C_c donné par la formule précédente

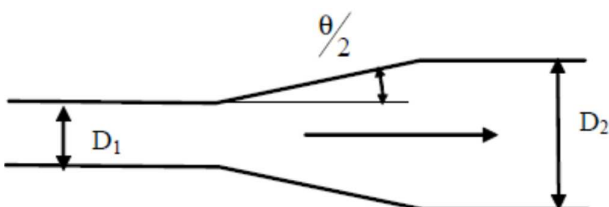
Convergent

Si le convergent est court et bien tracé, la perte de charge peut être négligée dans la plupart des applications

Divergent

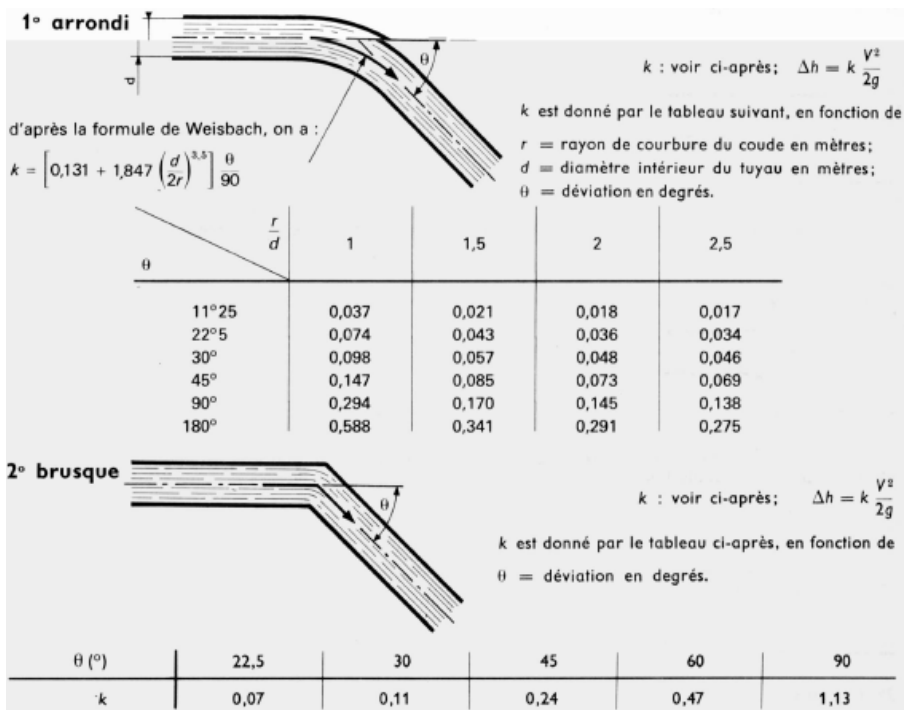
K peut varier notablement avec l'angle d'ouverture θ . K est approximé par:

$$K = 3,2 \cdot \left(\tan \frac{\theta}{2}\right)^2 \cdot \left(1 - \left(\frac{D_1}{D_2}\right)^2\right)^2$$



Appendix : Head losses/ Minor losses

Exemple des changements de direction



Coude arrondi

$$K = \left[0,13 + 1,847 \cdot \left(\frac{d}{2r} \right)^{3,5} \right] \cdot \frac{\theta}{90}$$

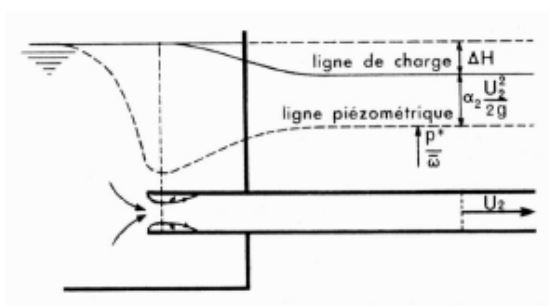
Coude à angle vif

$$K = \sin^2 \frac{\theta}{2} + \sin^4 \frac{\theta}{2}$$

35

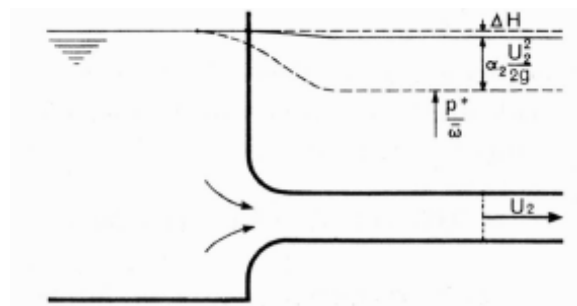
Appendix : Head losses/ Minor losses

Débouché d'une conduite dans un réservoir



Orifice de Borda

$$C_c \cong 0,5 \Rightarrow K \cong 1$$



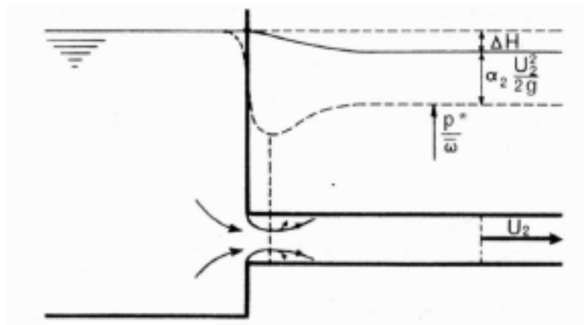
Orifice à bords arrondis et polis

$$C_c \cong 0,99 \Rightarrow K \cong 0$$

36

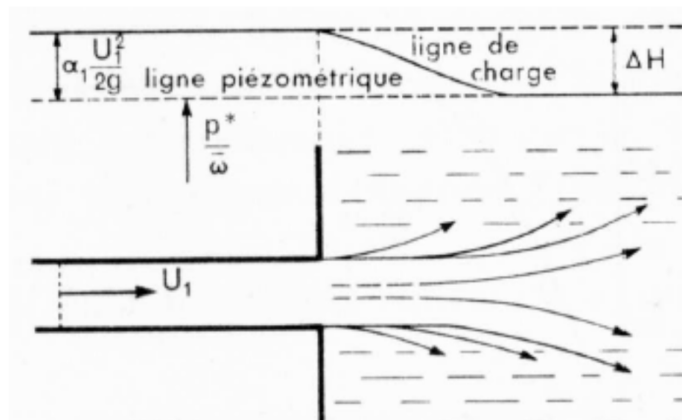
Appendix: Head losses/ Minor losses

Débouché d'une conduite dans un réservoir



Orifice à bords vifs

$$C_c \approx 0,6 \Rightarrow K \approx 0,5$$



Arrivée dans un réservoir

Il s'agit du cas particulier de **l'élargissement brusque** traité avec la formule de Borda $\Rightarrow K = 1$

La ligne piézométrique se raccorde avec la surface libre du réservoir.

37

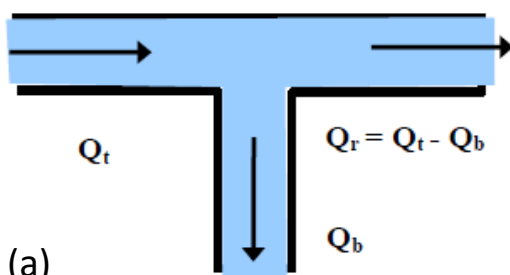
Appendix : Head losses/ Minor losses

Branchements à 90°

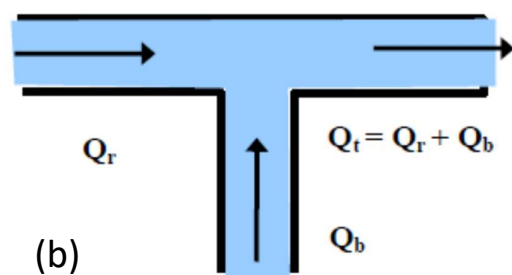
Des coefficients de perte de charge K_r et K_b sont respectivement affectés à la conduite rectiligne (indice r) et à la conduite de branchement (indice b) en fonction du rapport Q_b / Q_t . **Dans tous les cas la perte de charge est ensuite calculée avec la vitesse totale V_t .**

$$\Delta H_r = K_r \cdot \frac{V_t^2}{2g}$$

$$\Delta H_b = K_b \cdot \frac{V_t^2}{2g}$$



(a)



(b)

cas (a)	Q_b / Q_t	0	0,2	0,4	0,6	0,8	1
	K_r	0	0,02	0,06	0,07 à 0,15	0,15 à 0,21	0,35 à 0,4
	K_b	0,95 à 1	0,88 à 1,01	0,89 à 1,05	0,95 à 1,15	1,10 à 1,32	1,28 à 1,45

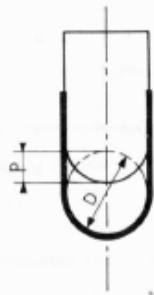
cas (b)	Q_b / Q_t	0	0,2	0,4	0,6	0,8	1
	K_r	0 à 0,04	0,17 à 0,27	0,30 à 0,46	0,41 à 0,57	0,51 à 0,60	0,55 à 0,60
	K_b			0,08 à 0,26	0,47 à 0,62	0,72 à 0,94	0,91 à 1,20

38

Appendix : Head losses/ Minor losses

Exemple de robinets

Robinet-vanne

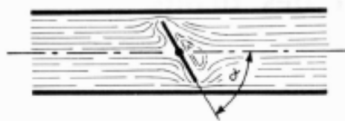


Le tableau suivant donne des valeurs expérimentales moyennes de k , en fonction de

p = distance de pénétration de l'obturateur dans la section, supposée circulaire, offerte par le robinet-vanne au passage du liquide, exprimée en mètres;

D = diamètre de cette section (diamètre intérieur du robinet-vanne), en mètres.

$\frac{p}{D}$	$\frac{1}{8}$	$\frac{2}{8}$	$\frac{3}{8}$	$\frac{4}{8}$	$\frac{5}{8}$	$\frac{6}{8}$	$\frac{7}{8}$
k	0,07	0,26	0,81	2,1	5,5	17	98



k : voir ci-après; $\Delta h = k \frac{V^2}{2g}$

Le tableau suivant donne des valeurs expérimentales moyennes de k , en fonction de α = angle formé par le papillon et l'axe de la conduite, en degrés.

α	5	10	15	20	30	40	45	50	60	70
k	0,24	0,52	0,90	1,5	3,9	11	19	33	120	750

Robinet à papillon

39

Appendix : Head losses/ Minor losses

Exemple de robinets

3° Robinets à tournant



k : voir ci-après; $\Delta h = k \frac{V^2}{2g}$

Le tableau suivant donne des valeurs expérimentales moyennes de k , en fonction de α = angle formé par l'axe de la lumière du boisseau — supposée à section circulaire et de même diamètre que l'intérieur du robinet — et l'axe de la conduite, en degrés.

α	5	10	15	25	35	45	55	65
k	0,05	0,29	0,75	3,1	9,7	31	110	490

Robinet à tournant

40